



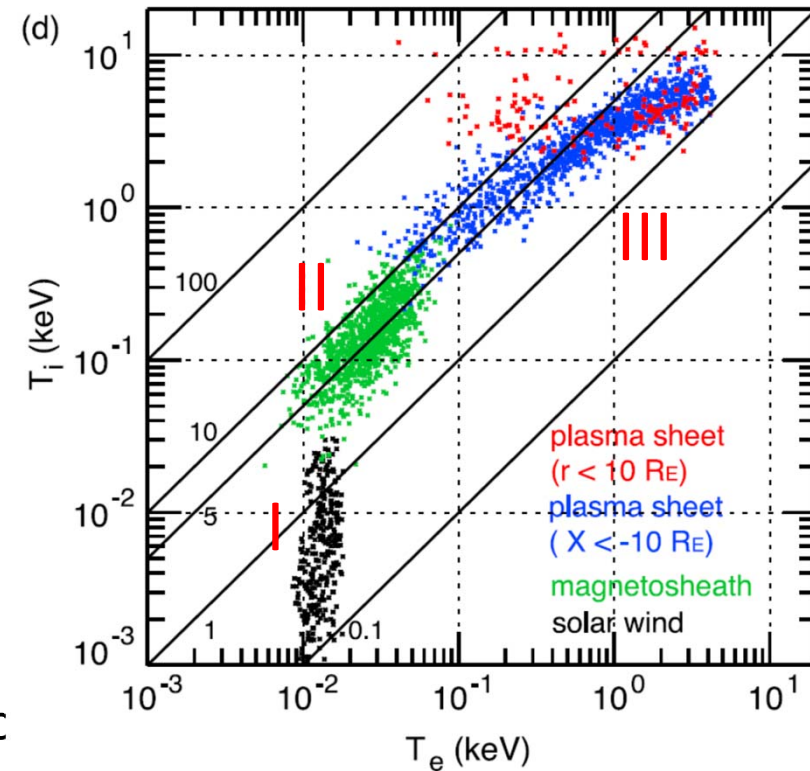
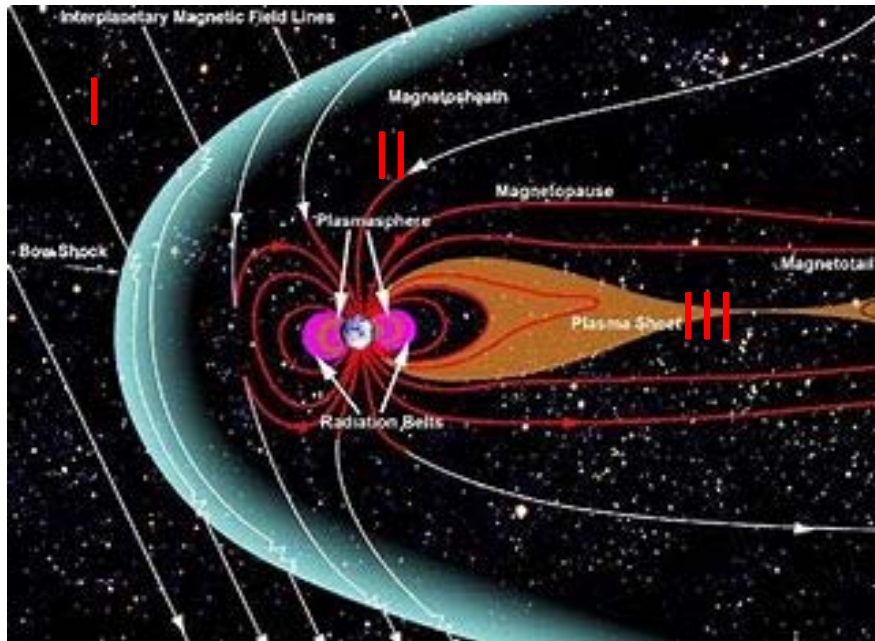
THE UNIVERSITY OF TOKYO

Energy partition of ions and electrons in the course of magnetic reconnection

Masahiro Hoshino
University of Tokyo

MH, ApJL (2018)

Observations of T_i/T_e in Magnetosphere



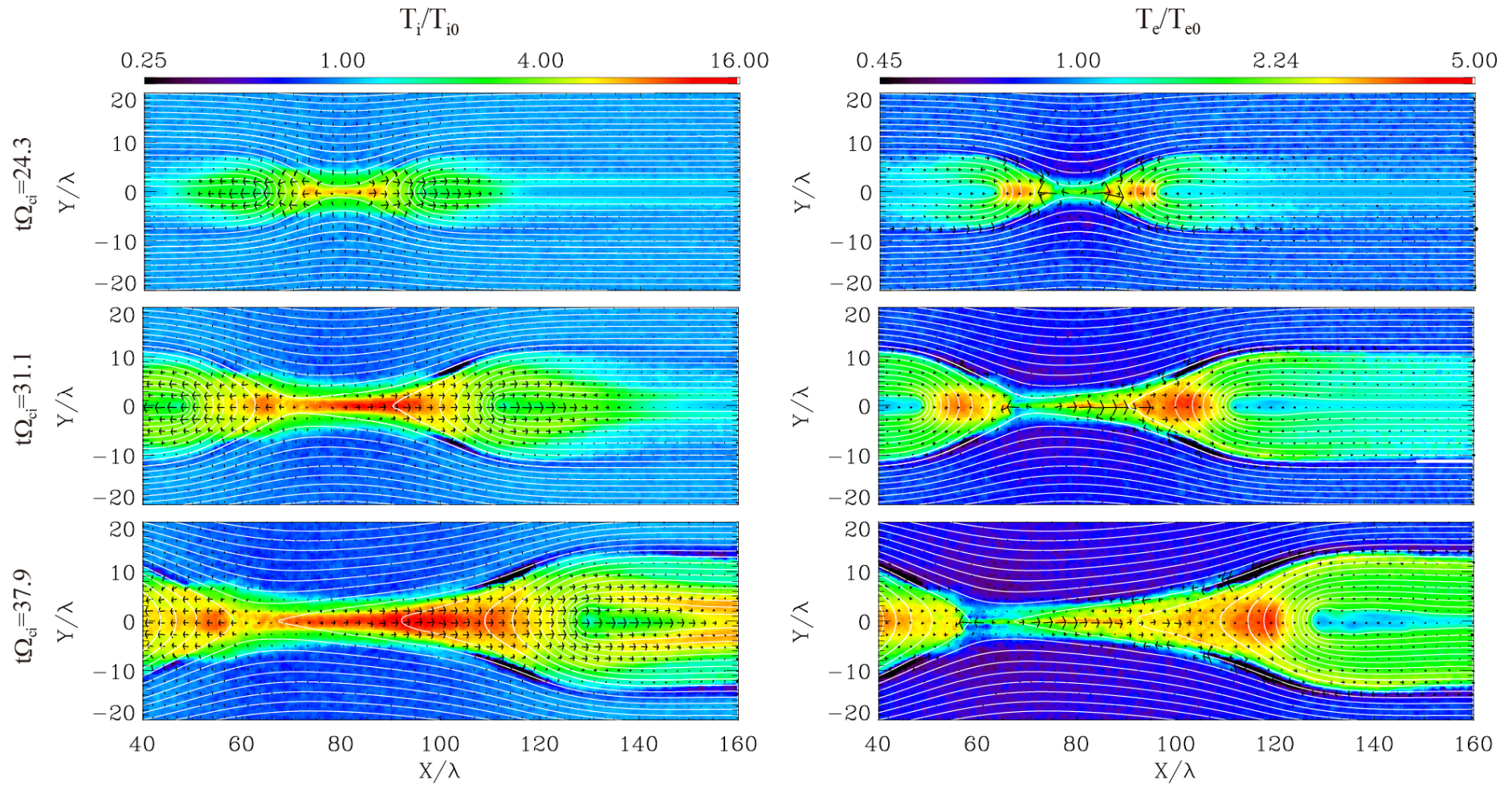
Hot ions in Earth's magnetotail are believed to be generated during magnetic reconnection

$$T_i/T_e = 5 \sim 10$$

(cf. Baumjohann+ JGR 1989; Eastwood+ PRL 2013; Phan+ GRL 2013)

Wang+ JGR 2012

T_i & T_e Heating in PIC simulation



cf. Simulation study of plasma heating: Wu+ PRL 2013; Shay+ PoP 2014; Haggerty+ GRL 2015

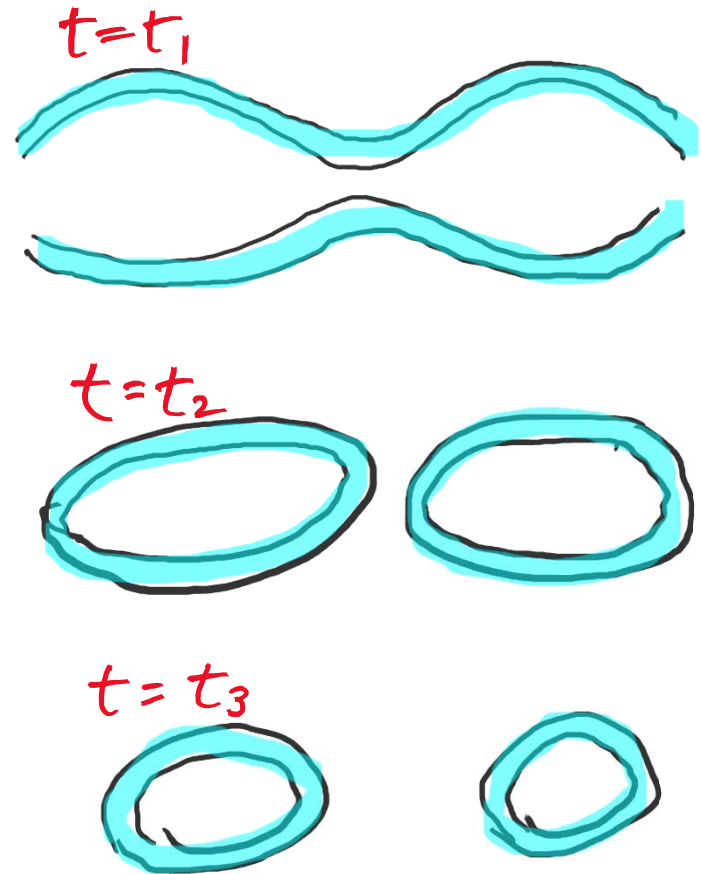
How can we understand the preferential ion heating?

Distinguish adiabatic heating and non-adiabatic heating

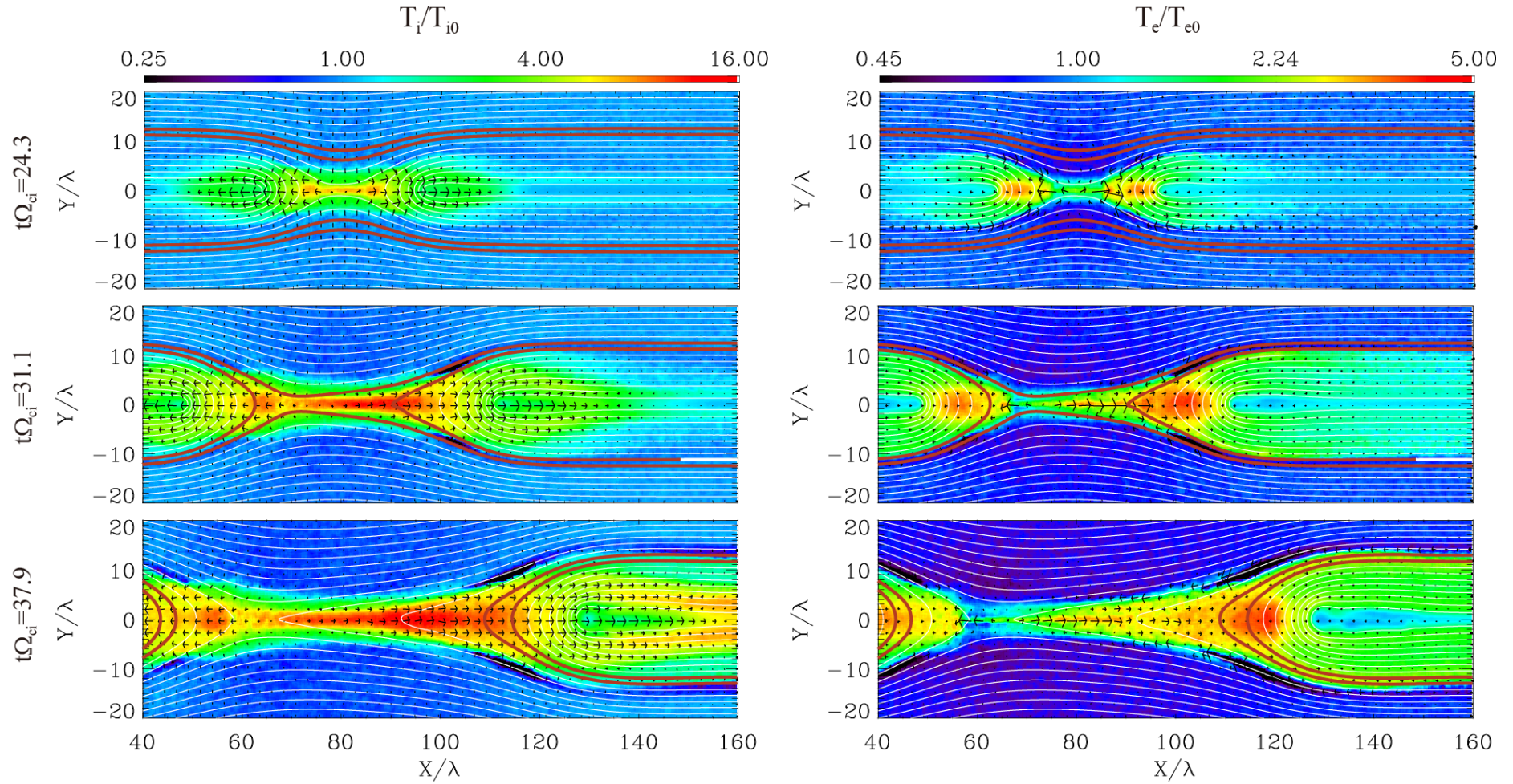
$$\frac{D}{Dt} \left(\frac{p}{\gamma - 1} \right) = \underbrace{\left(\frac{p}{\gamma - 1} \right) \frac{\gamma}{\rho} \frac{D\rho}{Dt}}_{\text{Adiabatic}} + \underbrace{Q_{heat}}_{\text{Nonadiabatic}}$$
$$Q_{heat} = \underbrace{\eta J^2}_{\text{Ohmic/Joule Heating}} + \underbrace{\text{others}}_{\text{Slow Shock, Turbulence...}}$$

If adiabatic heating, $PV^\gamma = const$

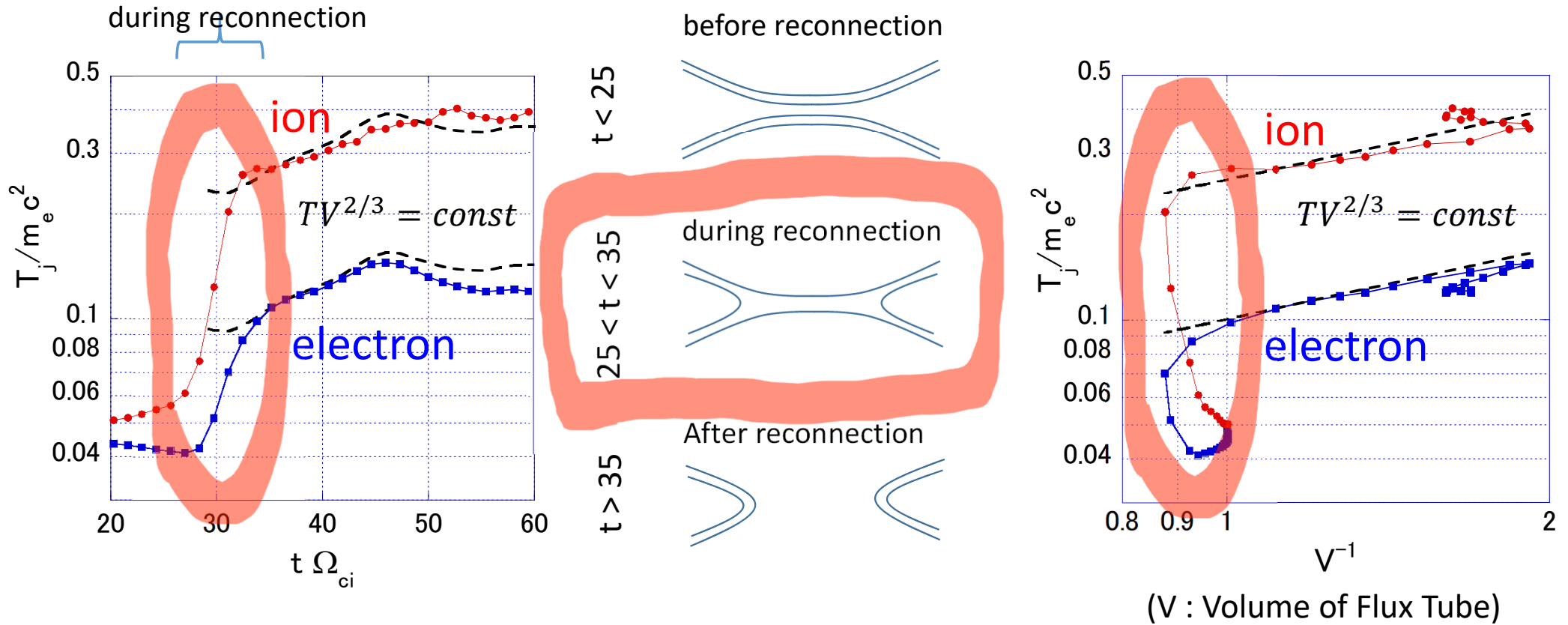
V is the volume of flux tube



Evolution of Magnetic Flux Tube

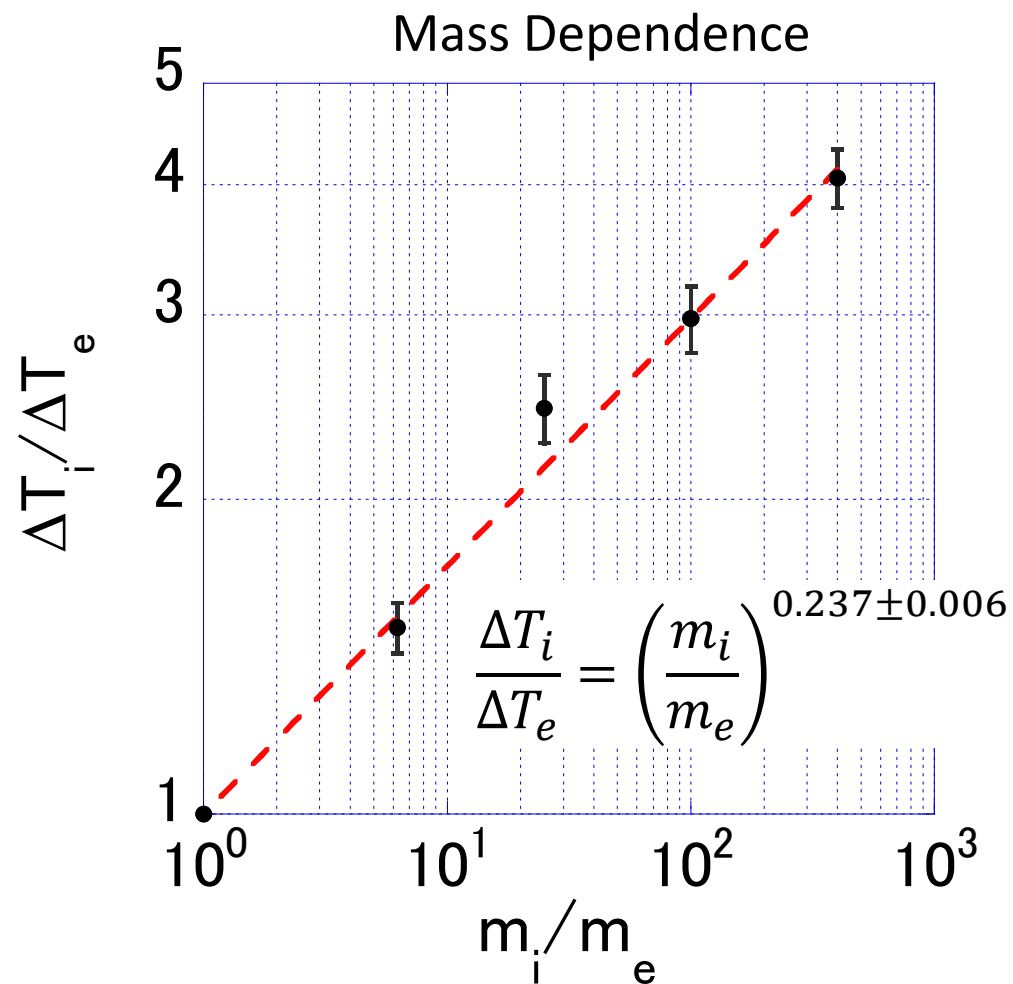
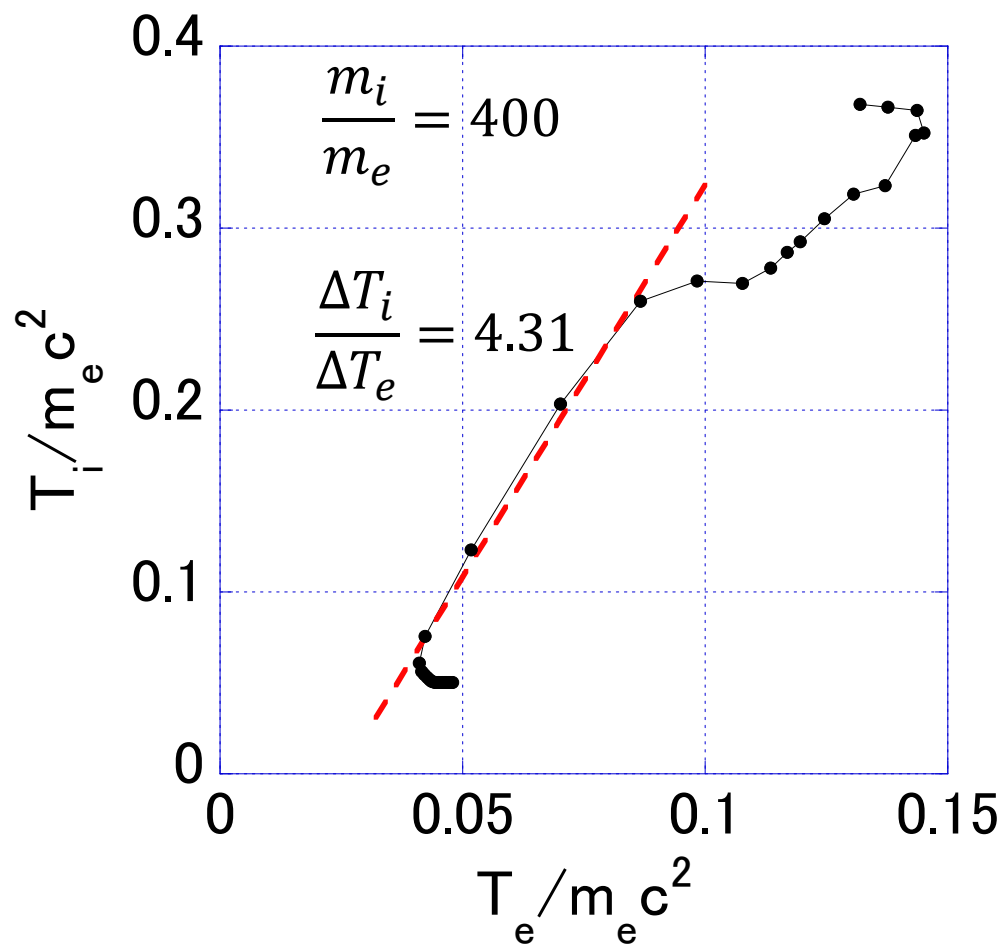


Time History & T-V Relation

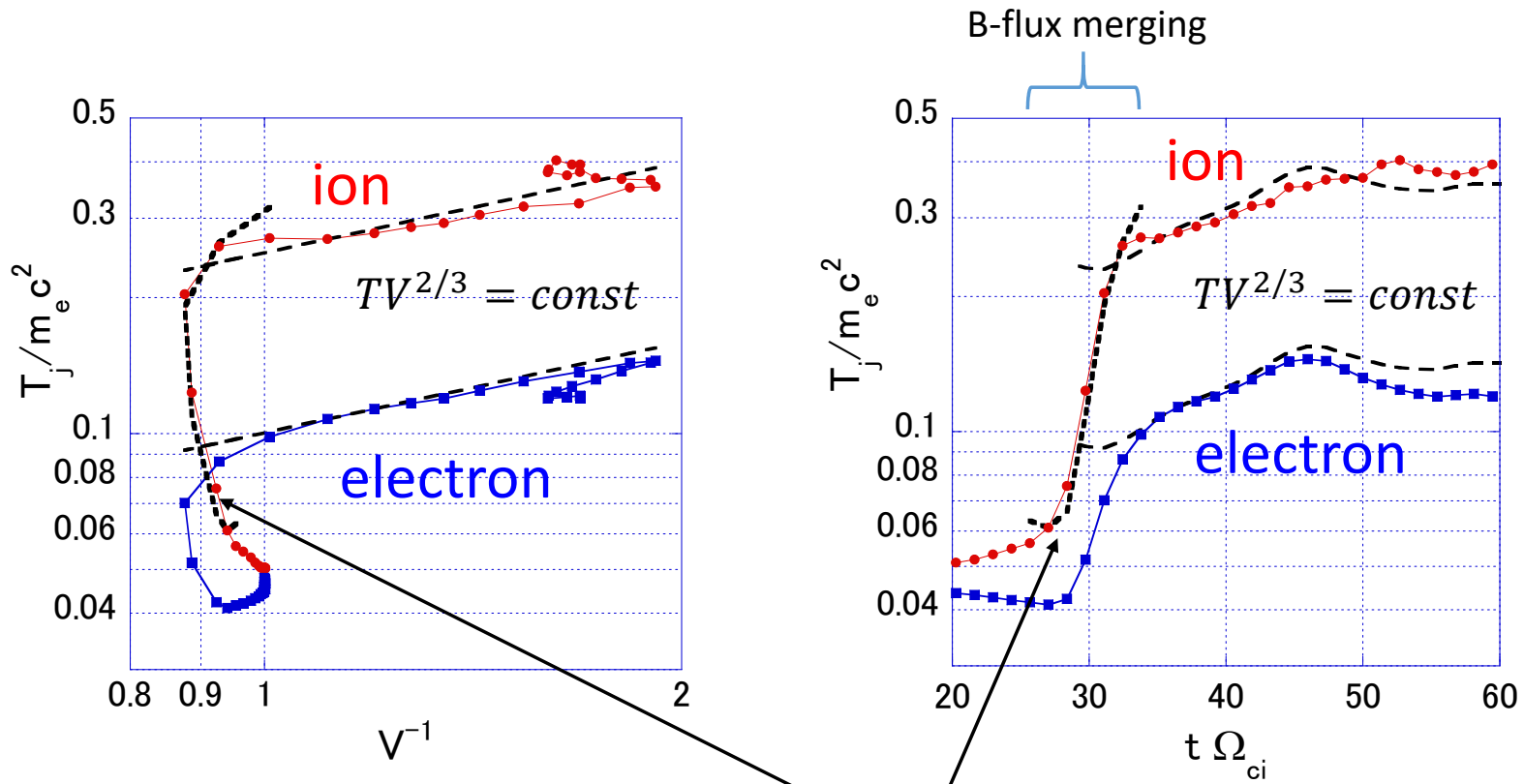


Adiabatic process after B-flux tube merging/re-connection

Ti/Te during B-flux merging stage



Thermodynamics of Reconnection

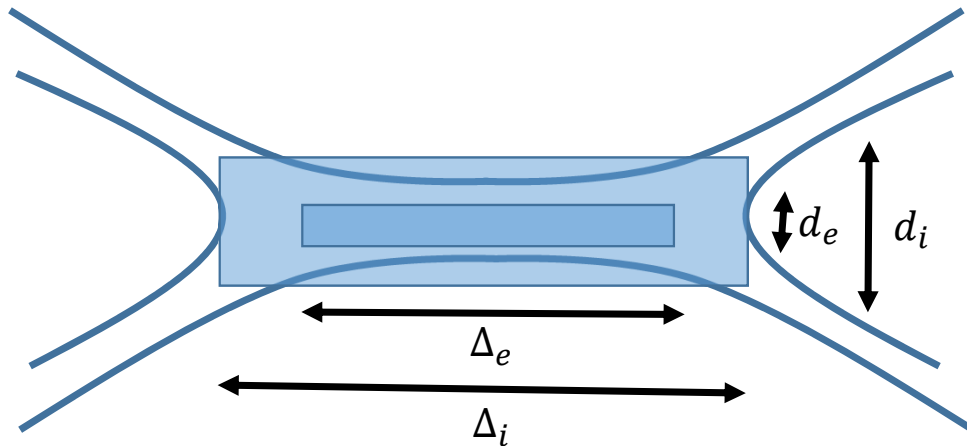


(V : Volume of Flux Tube)

Heating during
B-flux merging

$$\frac{\Delta T_i}{\Delta T_e} = \left(\frac{m_i}{m_e} \right)^{1/4}$$

Nonadiabatic heating during B-flux merging



collisionless/inertia conductivity

$$\sigma_j = \frac{ne^2}{m_j v_{c,j}} \propto \frac{\Delta_j}{(m_j T_j)^{1/2}}$$

(e.g. Coppi+ 1966; Hoh 1966; Galeev & Zeleny 1976)

meandering length

$$d_j = \sqrt{\frac{v_{th,j} \lambda}{\Omega_{cj}}} \propto (m_j T_j)^{1/4}$$

(e.g., Sonnerup 1971)

width of reconnection region

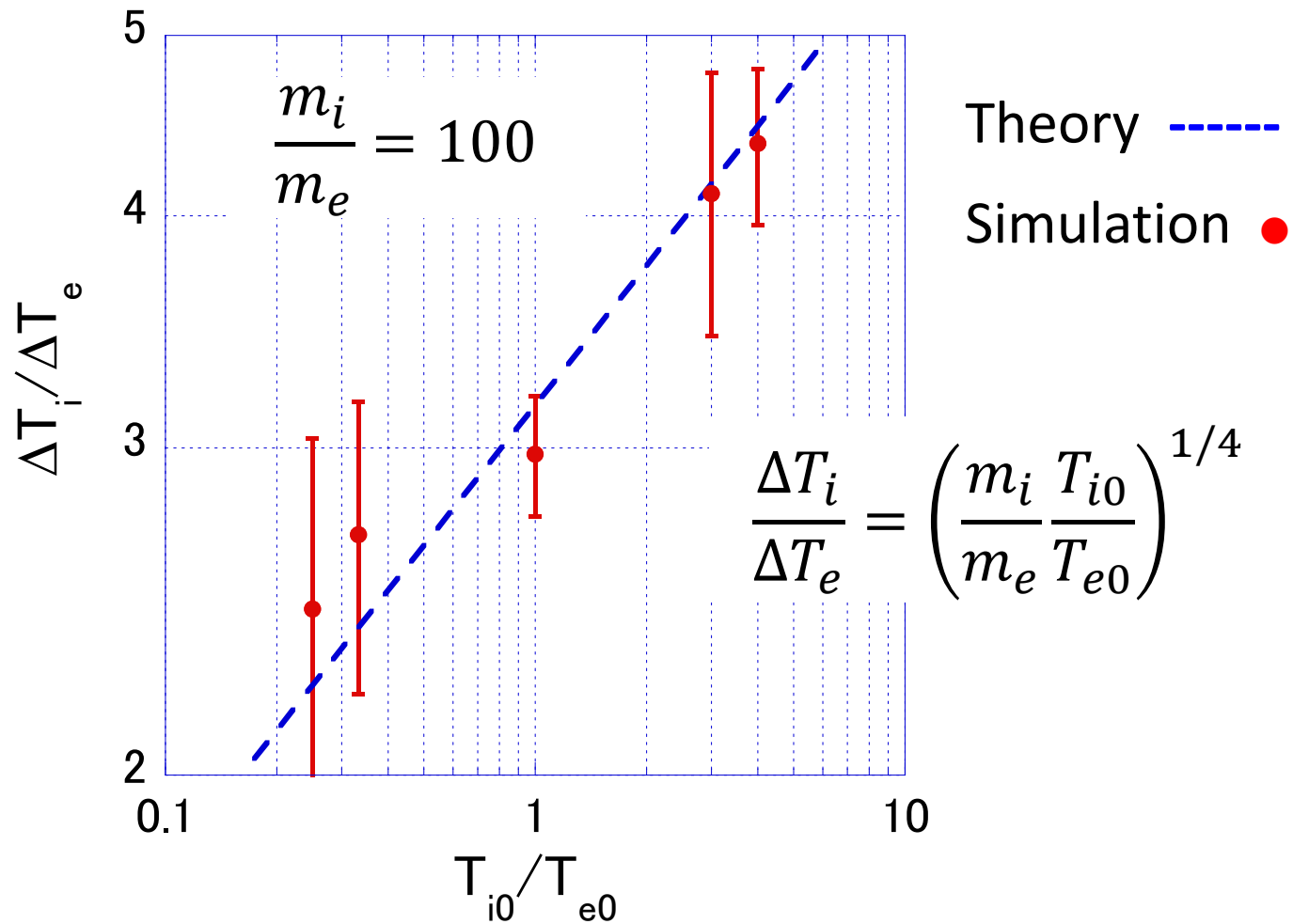
$$\Delta_j \propto (m_j T_j)^{1/4}$$

(e.g., Coroniti 1985)

$$\frac{\Delta T_i}{\Delta T_e} = \frac{\text{ion Joule heating}}{\text{ele. Joule heating}} = \frac{\sigma_i E^2 \Delta_i d_i}{\sigma_e E^2 \Delta_e d_e}$$

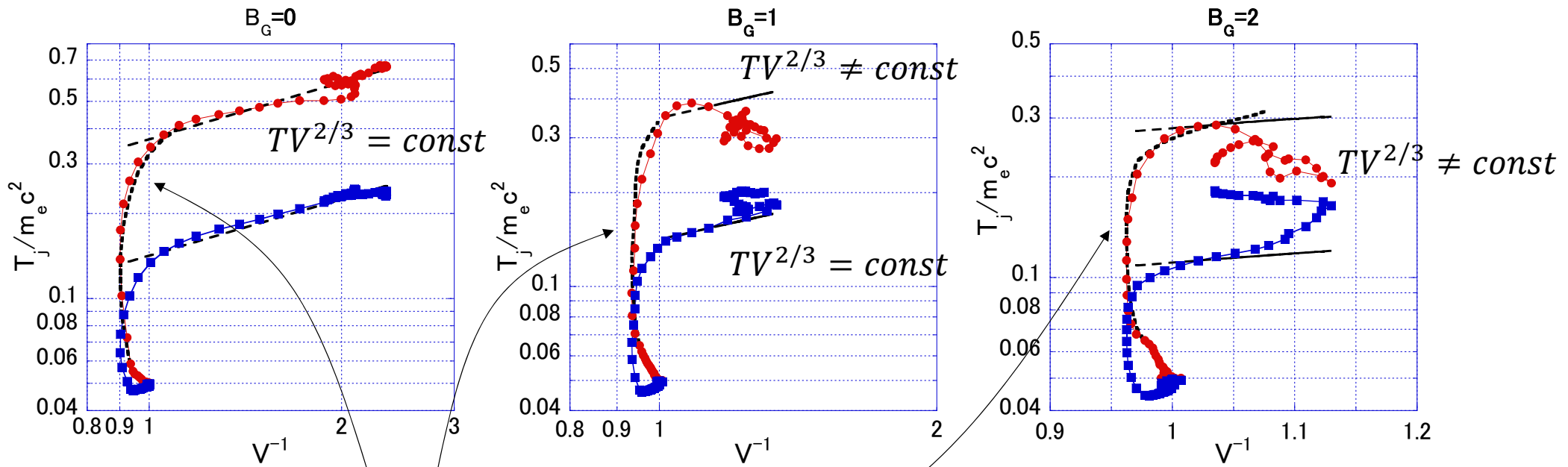
$$\frac{\Delta T_i}{\Delta T_e} = \frac{(m_i T_i)^{1/4}}{(m_e T_e)^{1/4}}$$

initial temperature dependence



Guide Magnetic Field Effect

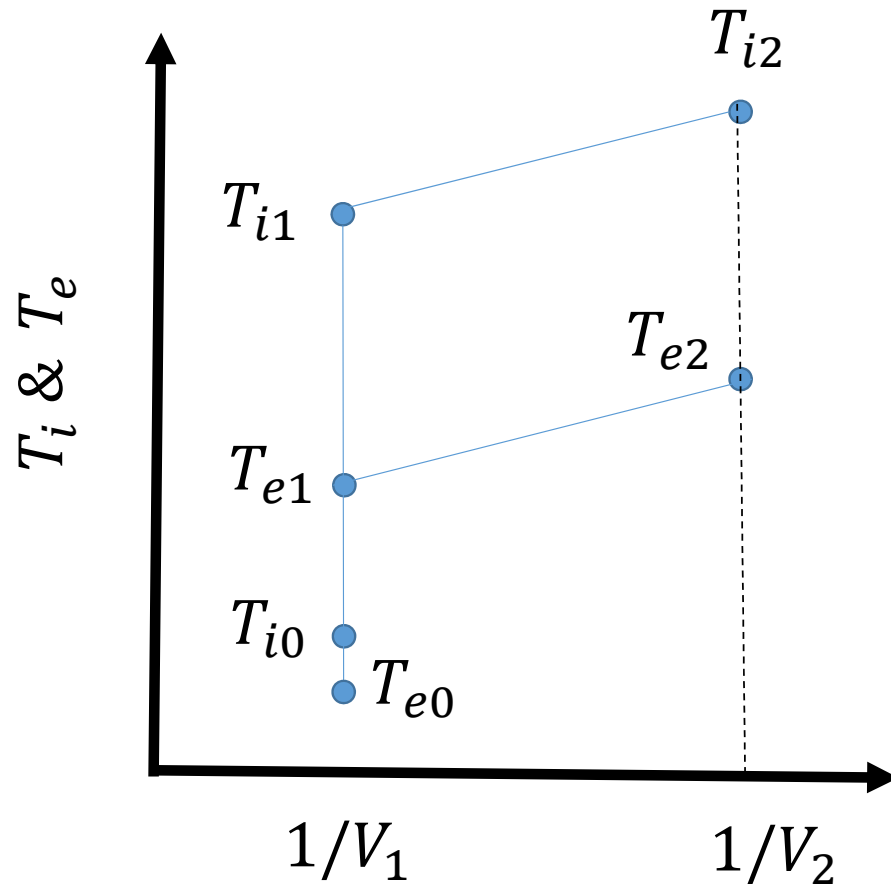
$$\frac{m_i}{m_e} = 100$$



$$\frac{\Delta T_i}{\Delta T_e} = \left(\frac{m_i}{m_e}\right)^{1/4}$$

weak compressibility

Thermodynamics of Reconnection



$$\frac{T_{i1} - T_{i0}}{T_{e1} - T_{e0}} = \left(\frac{m_i T_{i0}}{m_e T_{e0}} \right)^{1/4}$$

$$\frac{T_{e2}}{T_{e1}} = \left(\frac{V_1}{V_2} \right)^{\gamma-1} \quad \gamma = \frac{5}{3}$$

$$\frac{T_{i2}}{T_{i1}} = \left(\frac{V_1}{V_2} \right)^{\gamma-1}$$

Summary

- Energy Partition of Ion & Electron during Magnetic Reconnection

- Two distinct heating stages:

- Effective Ohmic/Joule Heating

$$\frac{\Delta T_i}{\Delta T_e} = \left(\frac{m_i T_{i0}}{m_e T_{e0}} \right)^{1/4}$$

- Adiabatic Compression (for anti-parallel reconnection)

$$\frac{D}{Dt} (TV^{\gamma-1}) = 0$$

- Ion heating is less effective with increasing B_G field