

# Horizon formation and far-from-equilibrium dynamics in strongly coupled plasma

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based on work with Paul Chesler: [arXiv:0906.4426](https://arxiv.org/abs/0906.4426), [arXiv:0906.4426](https://arxiv.org/abs/0906.4426)

# Thermal plasma physics from AdS/CFT

- Equilibrium ( $\mathcal{N} = 4$  SYM)

(static, Euclidean signature)

- equation of state
- correlation lengths, screening
- flavor physics
- finite volume
  - confinement/deconfinement
  - chemical potentials
  - rotation

SUGRA mode	$\mathfrak{g}_{R_y}^{CR_t}$	SYM operator	mass/ $\pi T$
$G_{00}$	$0_+^{++}$	$T_{00}$	2.3361
$a$	$0_-^{+-}$	$\text{tr } E \cdot B$	3.4041
$G_{ij}$	$2^{++}$	$T_{ij}$	3.4041
$\phi$	$0_+^{++}$	$\mathcal{L}$	3.4041
$G_{i0}$	$1^{+-}$	$T_{i0}$	4.3217
$B_{ij}$	$0_-^{+-}$	$\mathcal{O}_{ij}$	5.1085
$C_{ij}$	$0_+^{--}$	$\mathcal{O}_{30}$	5.1085
$B_{i0}$	$1^{--}$	$\mathcal{O}_{i0}$	6.6537
$C_{i0}$	$1^{-+}$	$\mathcal{O}_{3j}$	6.6537
$G_a^a$	$0_+^{++}$	$\text{tr } F^4$	7.4116

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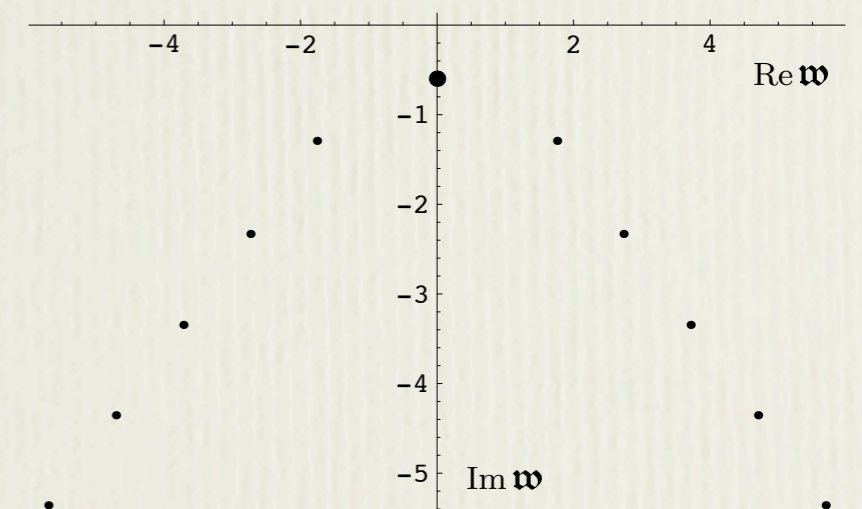
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- Near-equilibrium

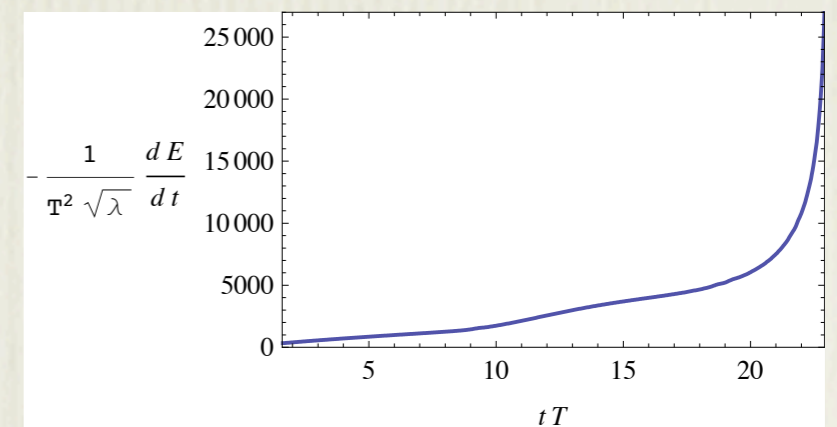
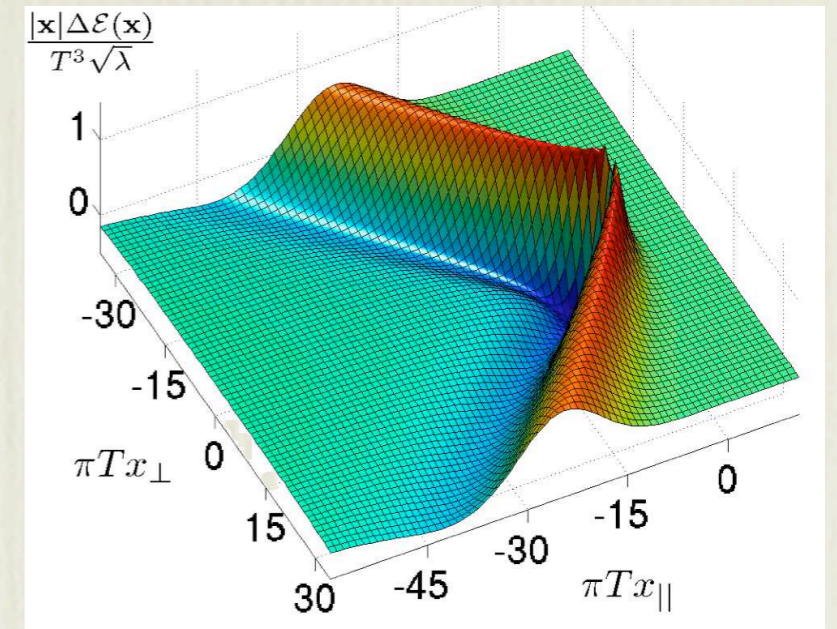
(real-time response, Minkowski signature)

- viscosity, diffusion
- quasi-normal modes, late time expansions
- photo-emission
- second-order transport coefficients
- non-linear conductivity



# Thermal plasma physics from AdS/CFT

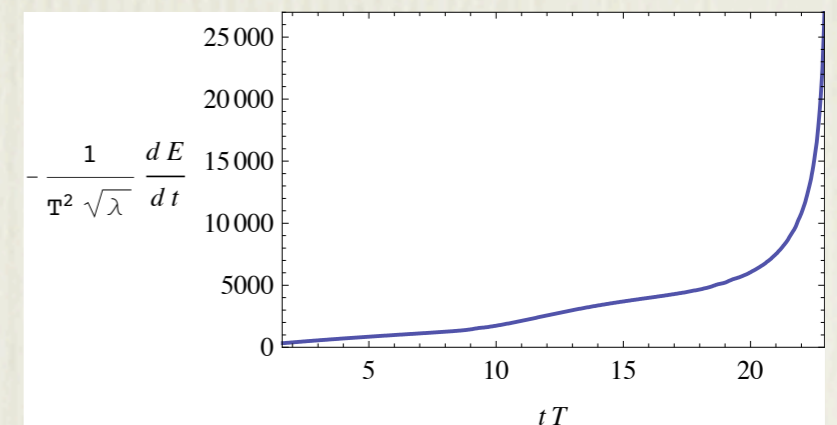
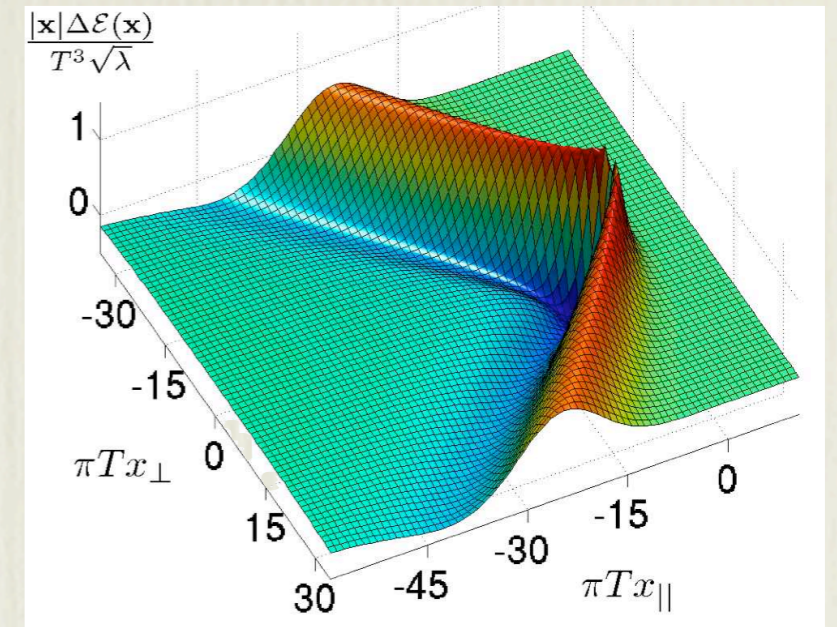
- Probe dynamics (classical string dynamics)
  - heavy quark drag
  - wakes, Brownian motion
  - heavy meson stability, dispersion
  - light quark jets



# Thermal plasma physics from AdS/CFT

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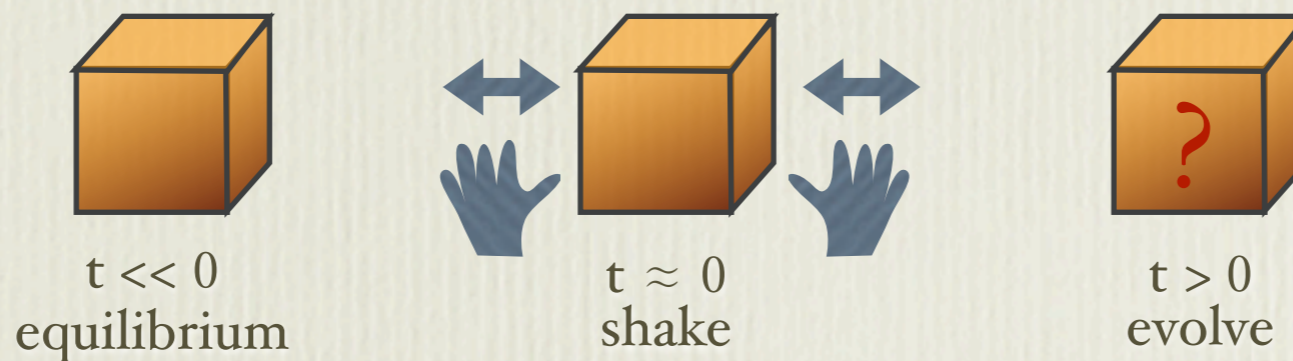
- Far-from-equilibrium dynamics ???

- plasma formation
- early thermalization
- turbulence

} initial value problems with non-trivial time-dependent bulk geometry

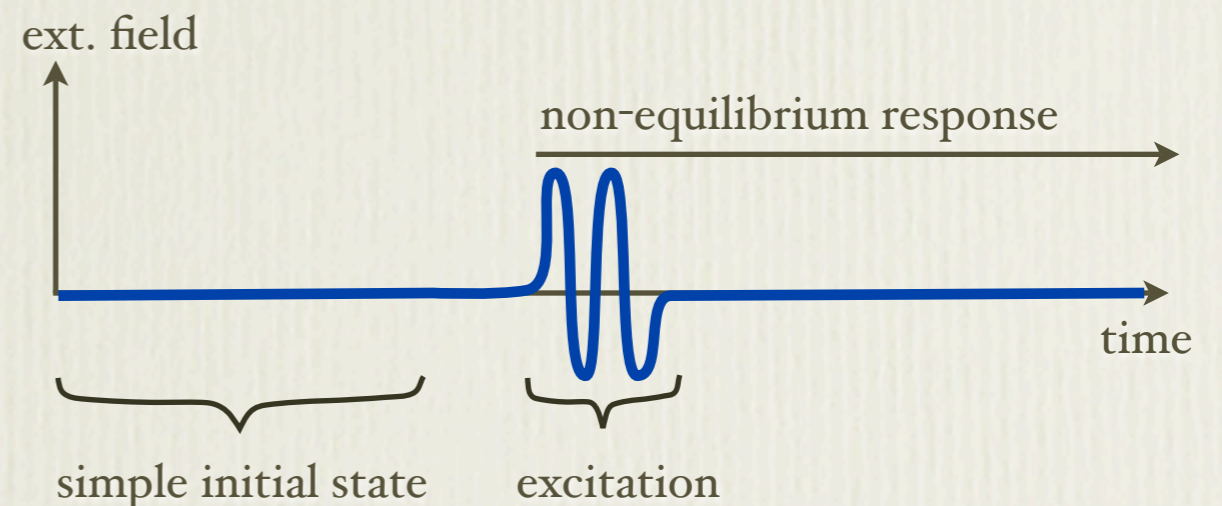
# Non-equilibrium initial states

- Specify complete density matrix  $\rho$  ? Ugh!
- Pick geometry on initial Cauchy surface ? Ugh!
- Want “operational” description:



∴ Specify time-dependent external fields

- ➔ time-dependent dynamics
- ➔ external work done on system



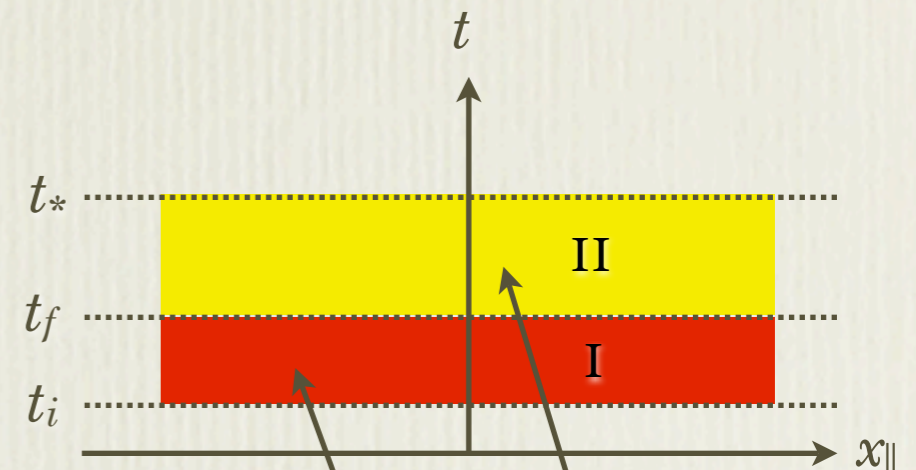
# Anisotropy dynamics

- Metric  $g^{\mu\nu}$  = external field coupling to stress-energy  $T^{\mu\nu}$

$\therefore$  time-dependent geometry  $\rightarrow$  non-equilibrium  $\langle T^{\mu\nu} \rangle$

- Case I: perfect spatial homogeneity, arbitrary anisotropy

$$ds^2 = -dt^2 + e^{f(t)}(dx^2 + dy^2) + e^{-2f(t)} dz^2$$



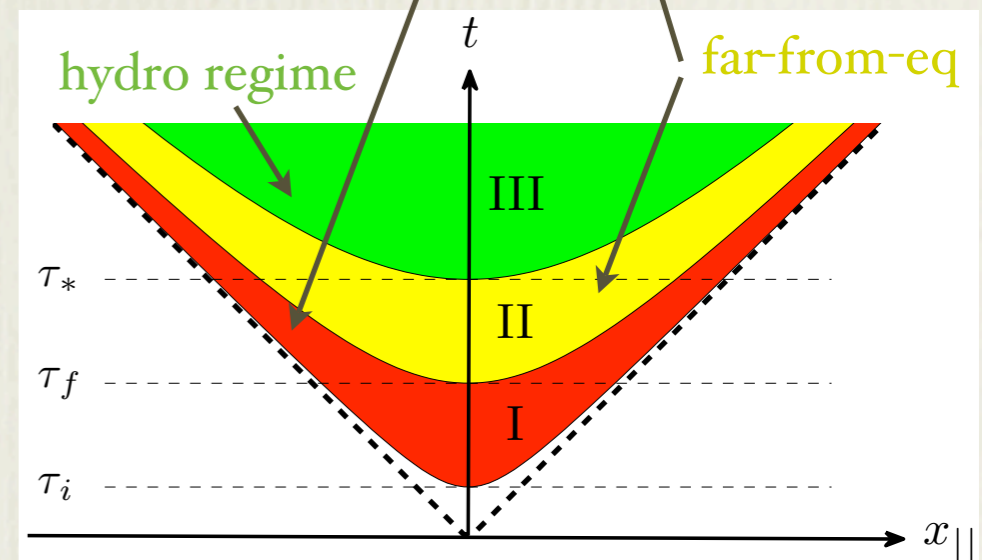
- Case II: boost invariance & transverse homogeneity

$$ds^2 = -d\tau^2 + e^{\gamma(\tau)} d\mathbf{x}_\perp^2 + \tau^2 e^{-2\gamma(\tau)} dy^2$$

proper time

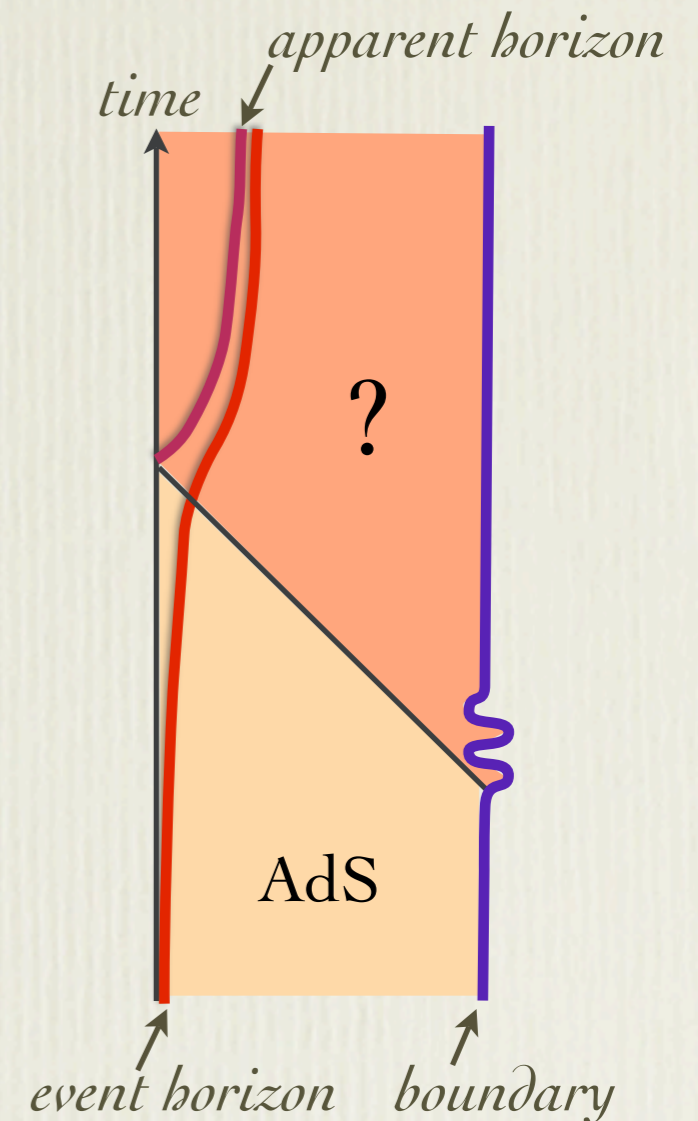
rapidity

$$\langle T^{\mu\nu}(t, \mathbf{x}) \rangle = \begin{bmatrix} \varepsilon(t) & & & \\ & p_\perp(t) & & \\ & & p_\perp(t) & \\ & & & p_\parallel(t) \end{bmatrix}$$



# Gravitational description

- Solve 5- $d$  Einstein equations with time-dependent boundary condition  $G^{AB} \rightarrow g^{\mu\nu}$  and simple initial condition (AdS or AdS-BH)
- Extract  $\langle T^{\mu\nu} \rangle$  from sub-leading near-boundary asymptotics
- Note:
  - time-dependent boundary conditions produce dynamic event horizon
  - “Teleological” event horizon growth occurs outside causal future of boundary time dependence
- ➔ event horizon area (pulled back to boundary) *cannot* represent entropy in non-equilibrium setting





# Practical issues (I)

- Coordinate choice:

✗ **Bad:** Fefferman-Graham or similar  $(r, t, \mathbf{x})$

✓ **Good:** Incoming Eddington-Finkelstein

$$ds^2 = -A(v, r) dv^2 + 2 dv dr + \Sigma(v, r)^2 \left[ e^{B(v, r)} (dx^2 + dy^2) + e^{-2B(v, r)} dz^2 \right]$$

- $v = \text{const.}$  on incoming (radial) null geodesics
- $dv/dr = \frac{1}{2}A$  on outgoing (radial) null geodesics
  - $g' \equiv \partial_r g$  = directional derivative along incoming null geodesics,
  - $\dot{g} \equiv \partial_v g + \frac{1}{2}A \partial_r g$  = directional derivative along outgoing null geodesics

- Boundary conditions as  $r \rightarrow \infty$ :

- Case I:  $A \rightarrow r^2, \quad \Sigma \rightarrow r, \quad B \rightarrow f(v)$
- Case II:  $A \rightarrow r^2, \quad \Sigma \rightarrow r \tau^{1/3}, \quad B \rightarrow -2/3 \ln \tau + \gamma(\tau)$

# Einstein equations

- $R_{MN} - \frac{1}{2} G_{MN}(R + 2\Lambda) = 0$
- Non-trivial components:  $uv, rr, vr, zz, xx+yy$ 
  - ➔ 5 equations, 3 unknown functions  $(A, B, \Sigma)$
- Need to separate dynamics from constraints
  - ➔  $0 = \Sigma (\dot{\Sigma})' + 2\Sigma' \dot{\Sigma} - 2\Sigma^2,$
  - $0 = \Sigma (\dot{B})' + \frac{3}{2}(\Sigma' \dot{B} + B' \dot{\Sigma}),$
  - $0 = A'' + 3B' \dot{B} - 12\Sigma' \dot{\Sigma}/\Sigma^2 + 4,$
  - $0 = \ddot{\Sigma} + \frac{1}{2}(\dot{B}^2 \Sigma - A' \dot{\Sigma}),$  ← boundary value constraint
  - $0 = \Sigma'' + \frac{1}{2}B'^2 \Sigma,$  ← initial value constraint
- N.B.:  $A =$  non-dynamical auxiliary field

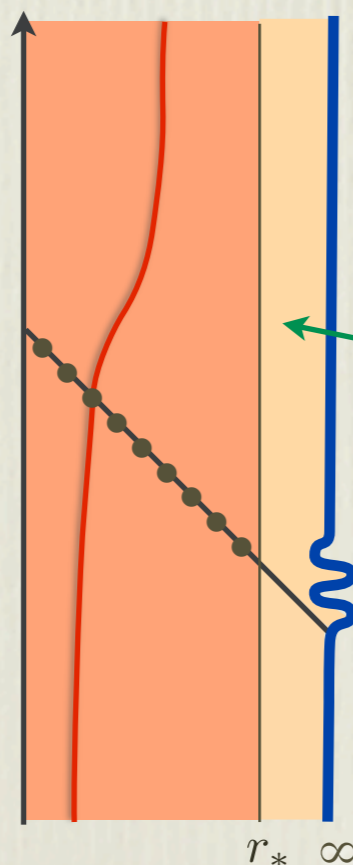
# Practical issues (II)

- Need to solve for “velocities,”  $\partial_v B$ ,  $\partial_v \Sigma$ , and auxiliary field  $A$

$$\dot{\Sigma}(r, v) = -\frac{2}{\Sigma(r, v)^2} \int_r dw \Sigma(w, v)^3$$

$$\dot{B}(r, v) = -\frac{3}{\Sigma(r, v)^{3/2}} \int_r dw \frac{B'(w, v)}{\Sigma(w, v)^{3/2}} \int_w d\bar{w} \Sigma(\bar{w}, v)^3$$

- Discretize  $r \rightarrow \infty$  system of coupled ODEs
- Must treat near-boundary behavior accurately
  - ➔ match discretized numerics to large  $r$  asymptotics



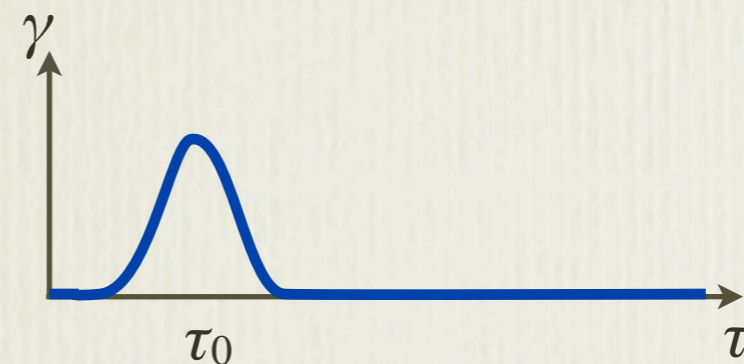
$$\left\{ \begin{aligned} A(r, v) &= \sum_{n=0} [a_n(v) + \alpha_n(v) \log r] r^{2-n}, \\ B(r, v) &= \sum_{n=0} [b_n(v) + \beta_n(v) \log r] r^{-n}, \\ \Sigma(r, v) &= \sum_{n=0} [s_n(v) + \sigma_n(v) \log r] r^{1-n}. \end{aligned} \right.$$

# Practical issues (III)

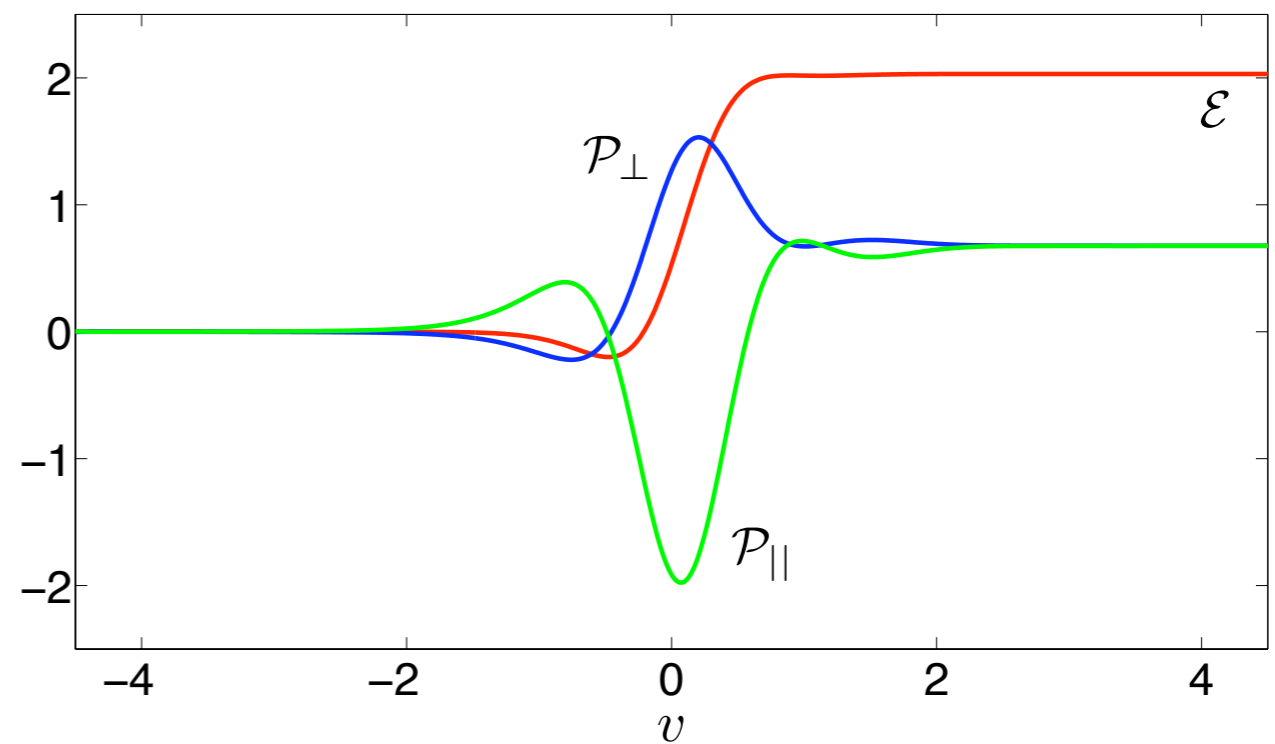
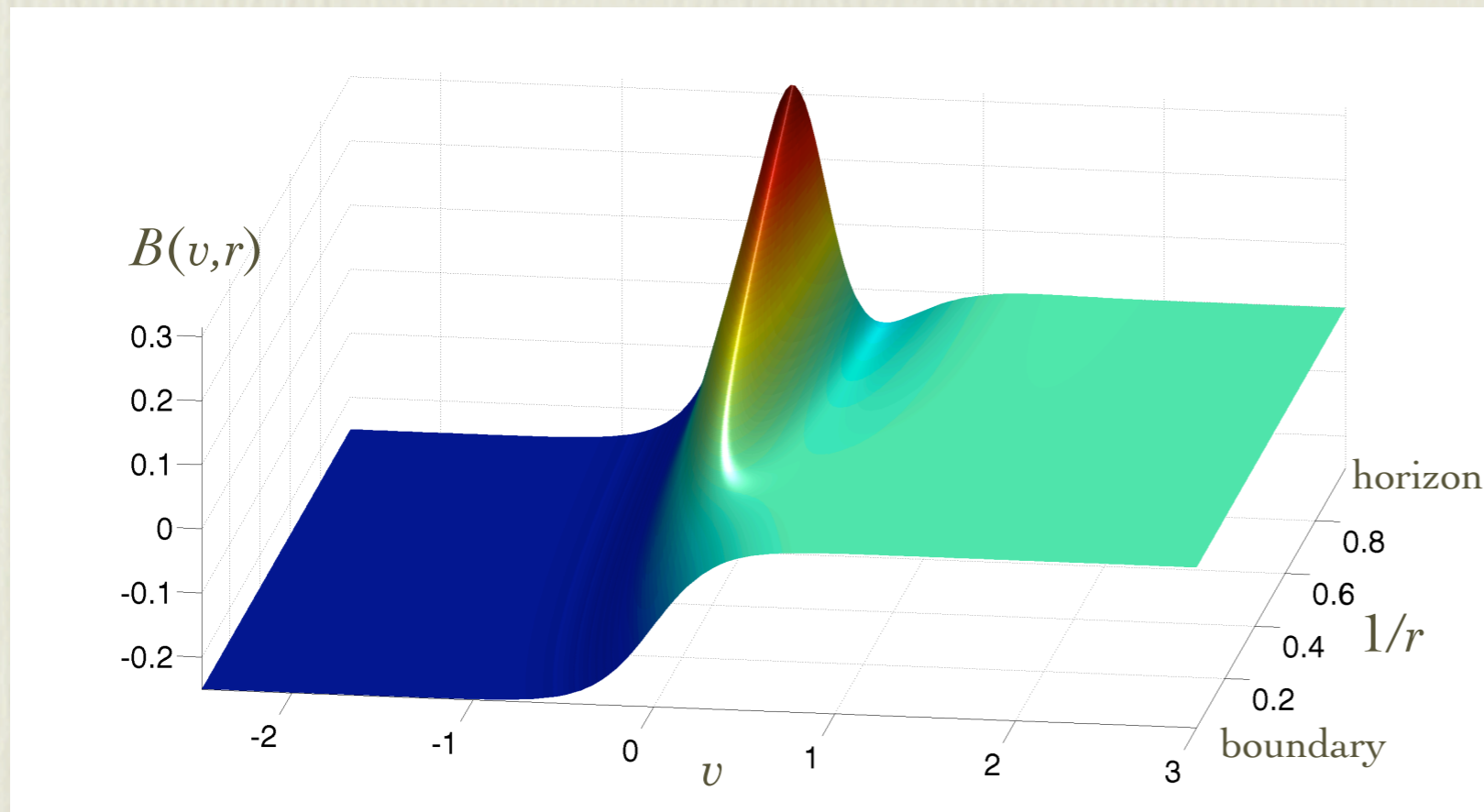
- Must remove residual reparameterize freedom:  $r \rightarrow r + \alpha(v)$ 
  - ✗ **Bad:** fix coordinate location of event horizon
  - ✓ **Good:** fix  $a_1 = 0$
- Must excise region surrounding singularity:  $r < r_{\min}(v) < r_{\text{horizon}}(v)$
- Must choose specific boundary time dependence

$$\text{Ex: } f(v) = \frac{1}{2}c [1 - \tanh(v/\tau)] \quad \gamma(\tau) = ch(\tau - \tau_0)^6 e^{-1/h(\tau - \tau_0)}$$

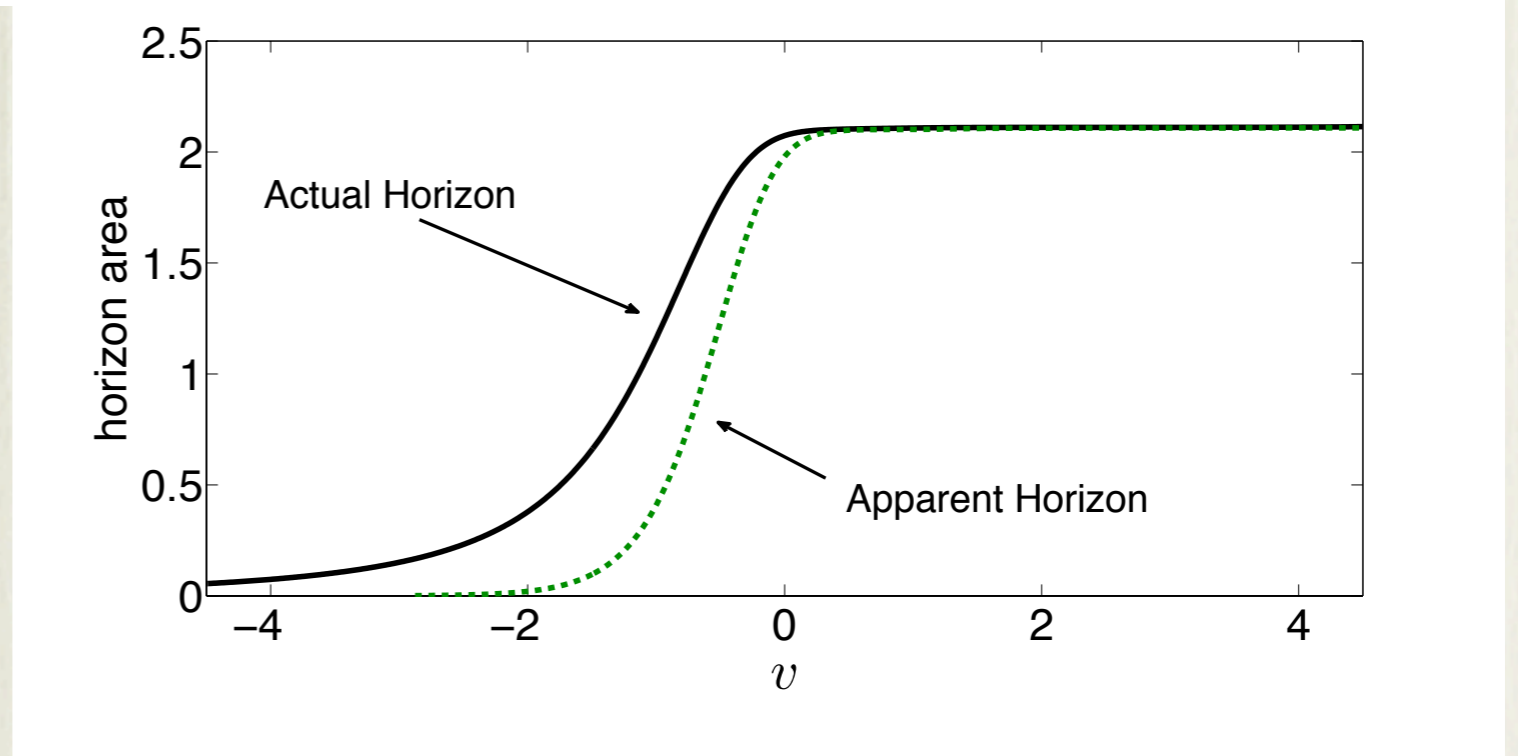
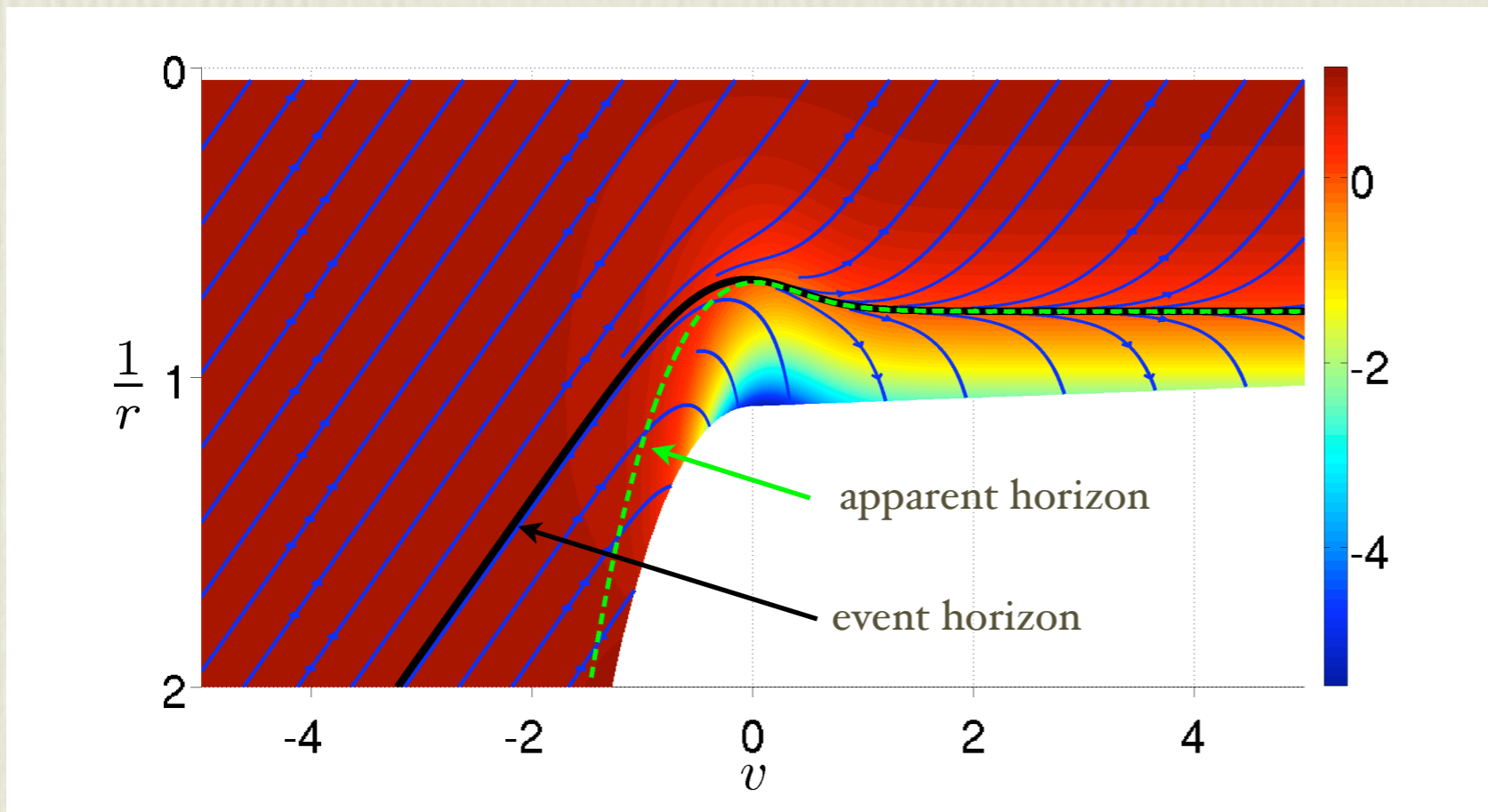
$$h(\delta\tau) = 1 - (\delta\tau)^2/\Delta^2$$



# Case I: Results



# Case I: Horizon area



# Case I: Isotropization time

$ c $	1	1.5	2	2.5	3	3.5	4
$\tau T$	0.23	0.31	0.41	0.52	0.65	0.79	0.94
$\tau_{\text{iso}} T$	0.67	0.68	0.71	0.92	1.2	1.5	1.8
$\tau_{\text{iso}}/\tau$	3.0	2.2	1.7	1.8	1.8	1.9	1.9

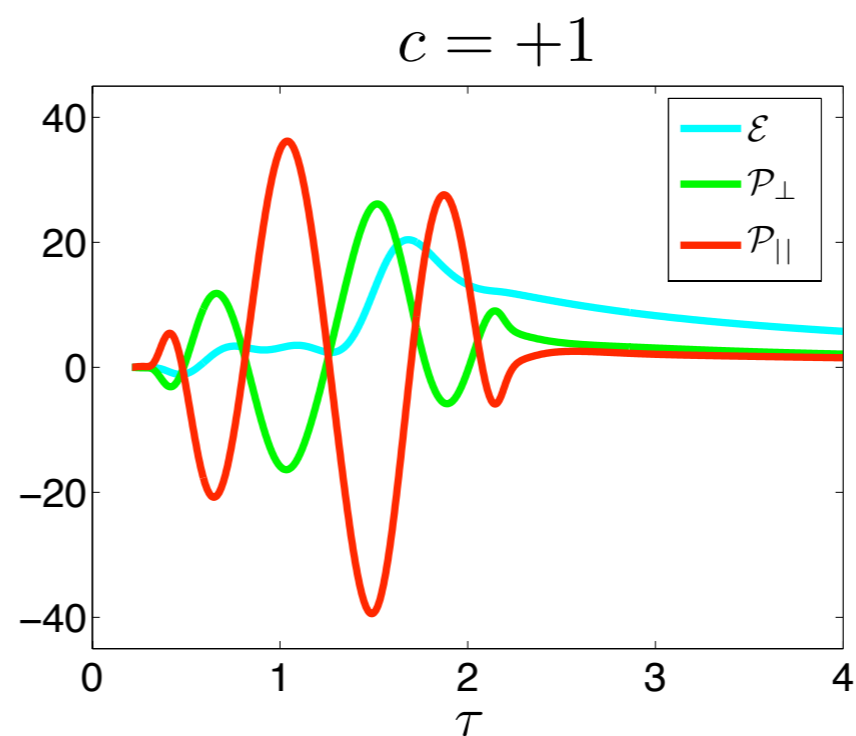
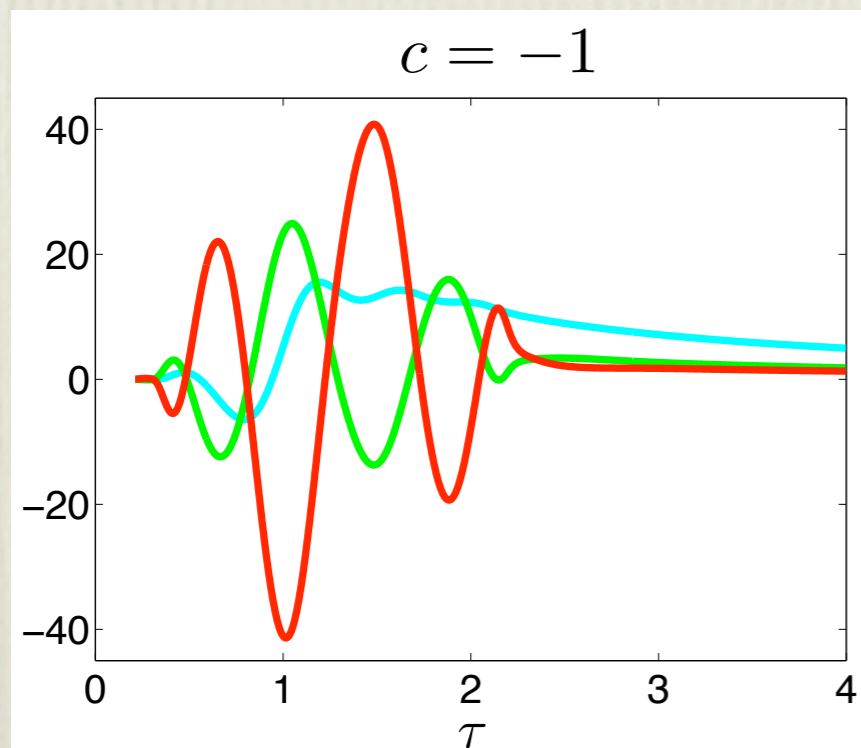
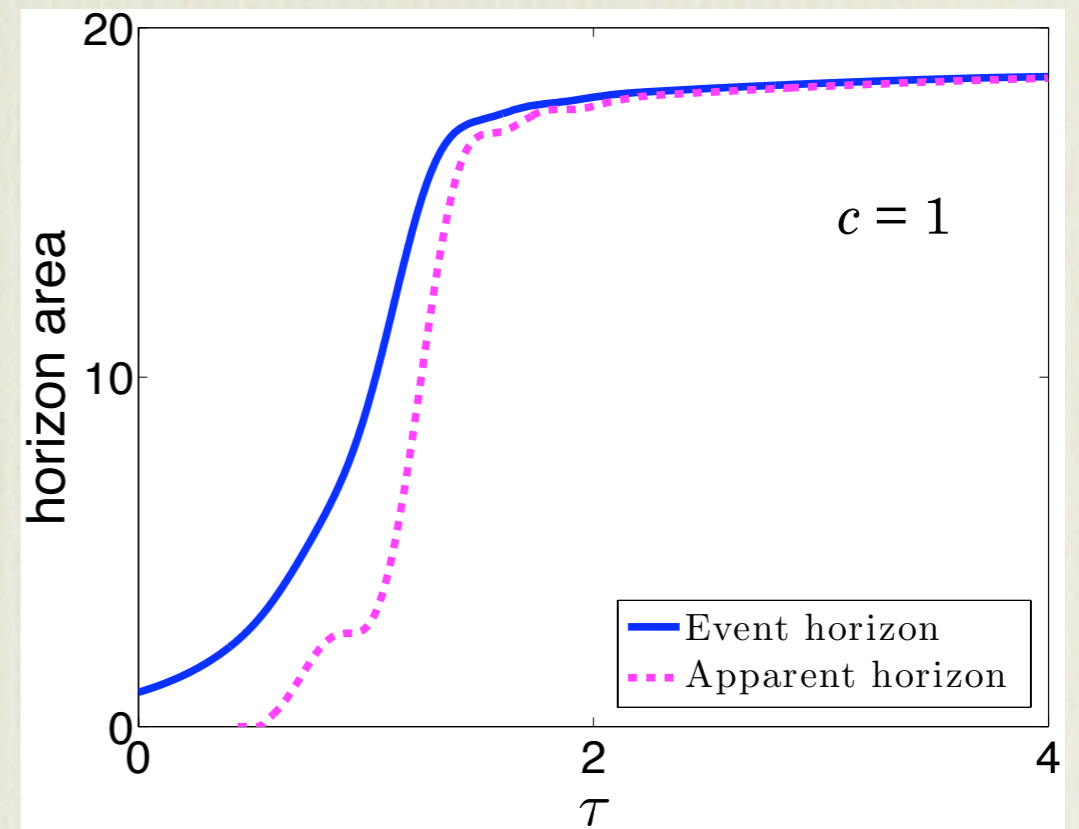
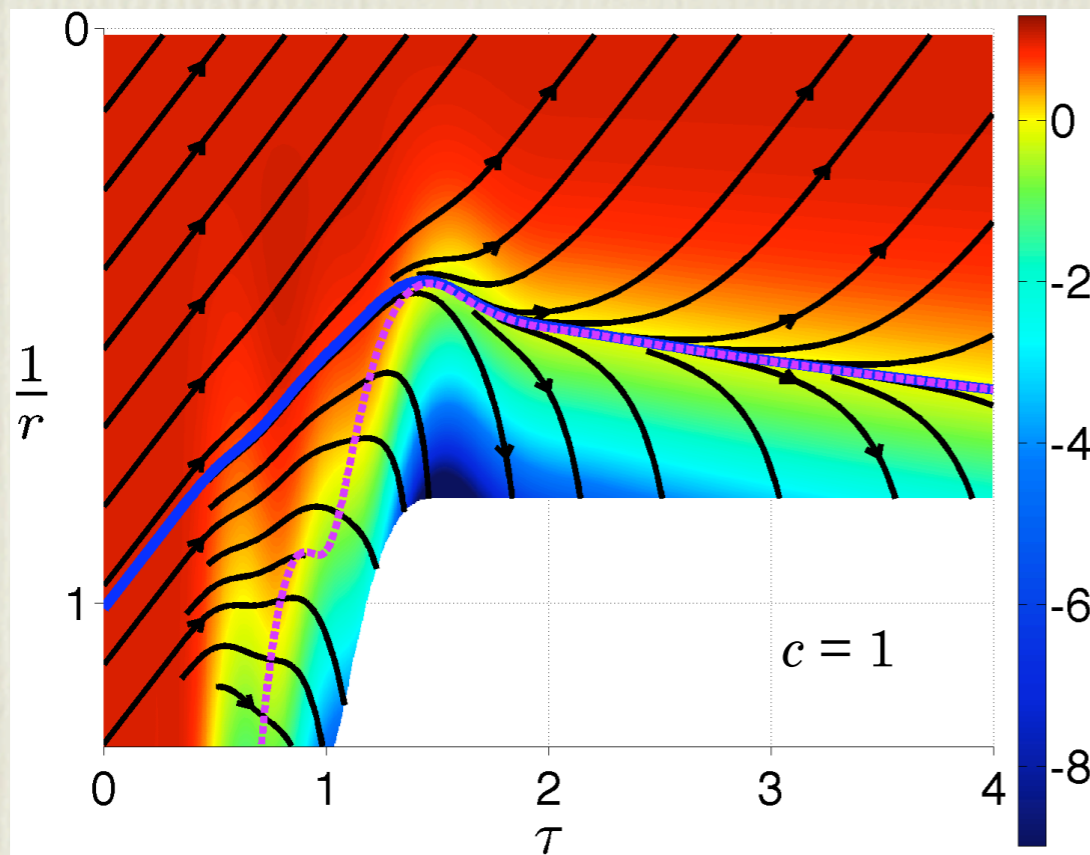
$T$  = final equilibrium temperature

$\tau_{\text{iso}}$  = isotropization time

$\tau$  = plasma creation time scale

$\tau_{\text{iso}} \approx 0.7/T \Rightarrow \tau_{\text{iso}} \approx 0.5 \text{ fm}/c$  at  $T \approx 350 \text{ MeV}$  --- relevant at RHIC???

# Case II: Results

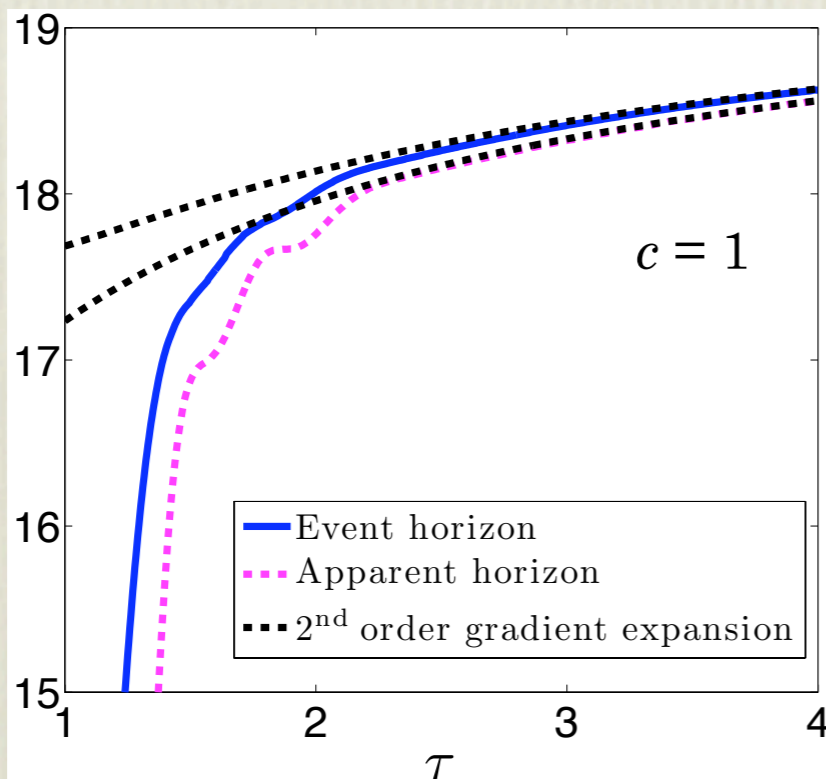
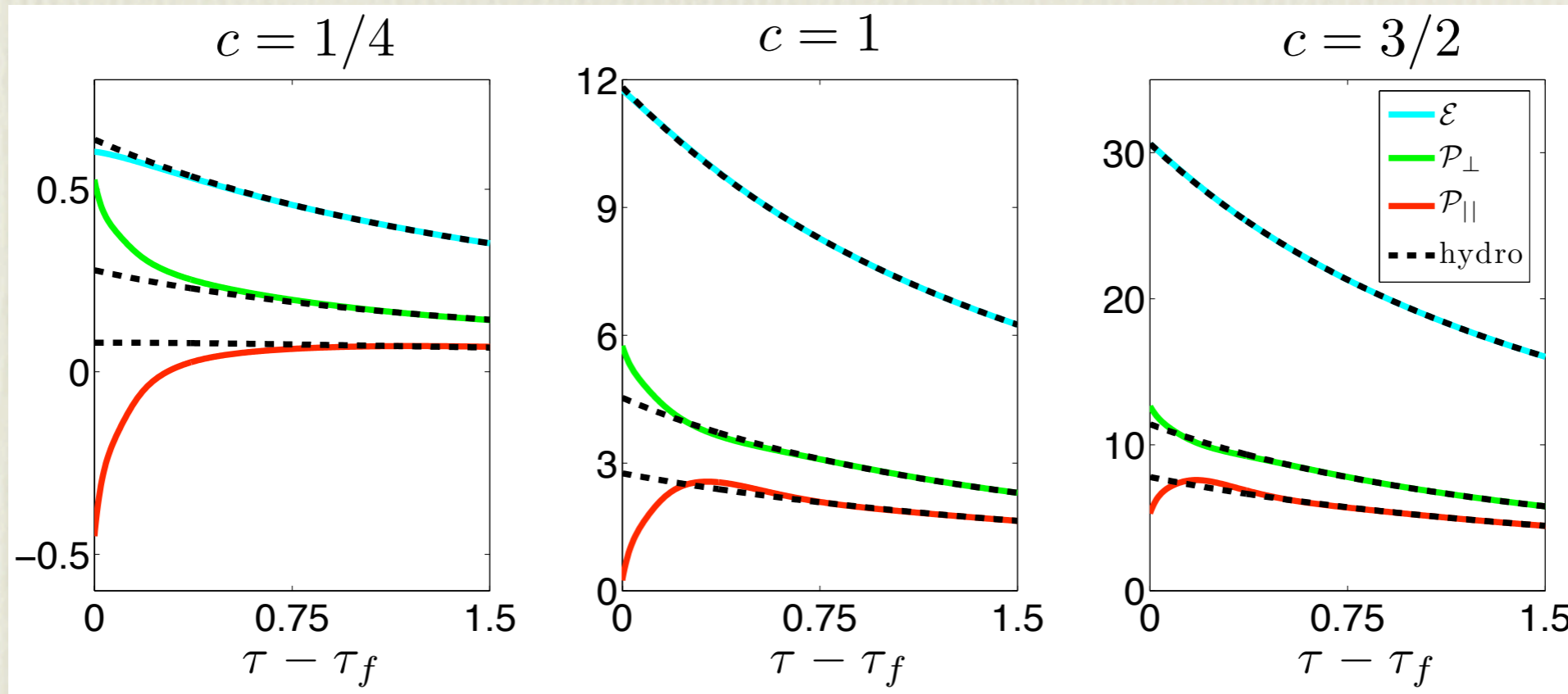


$$\tau_i = 0.25$$

$$\tau_f = 2.25$$



# Case II: Hydro comparison



$$\mathcal{E} = \frac{3\pi^4\Lambda^4}{4(\Lambda\tau)^{4/3}} \left[ 1 - \frac{2C_1}{(\Lambda\tau)^{2/3}} + \frac{C_2}{(\Lambda\tau)^{4/3}} \right], \quad C_1 = \frac{1}{3\pi}$$

$$\mathcal{P}_\perp = \frac{\pi^4\Lambda^4}{4(\Lambda\tau)^{4/3}} \left[ 1 - \frac{C_2}{3(\Lambda\tau)^{4/3}} \right], \quad C_2 = \frac{2 + \ln 2}{18\pi^2}$$

$$\mathcal{P}_\parallel = \frac{\pi^4\Lambda^4}{4(\Lambda\tau)^{4/3}} \left[ 1 - \frac{2C_1}{(\Lambda\tau)^{2/3}} + \frac{5C_2}{3(\Lambda\tau)^{4/3}} \right]$$

$$A_{\text{EH}} = \pi^3\Lambda^2 \left[ 1 - \frac{1}{2\pi(\Lambda\tau)^{2/3}} + \frac{6 + \pi + 6\ln 2}{24\pi^2(\Lambda\tau)^{4/3}} \right]$$

$$A_{\text{AH}} = \pi^3\Lambda^2 \left[ 1 - \frac{1}{2\pi(\Lambda\tau)^{2/3}} + \frac{2 + \pi + \ln 2}{24\pi^2(\Lambda\tau)^{4/3}} \right],$$

# Case II: Relaxation time

$c$	-2	-3/2	-1	-1/2	-1/4	1/4	1/2	1	3/2	2
$\tau_*$	2.2	2.3	2.4	2.7	3.1	3.1	2.7	2.4	2.3	2.2
$T_*$	0.93	0.77	0.60	0.40	0.27	0.27	0.41	0.62	0.80	0.97
$\Lambda\tau_*$	3.1	2.5	1.9	1.2	0.87	0.89	1.3	1.9	2.6	3.3
$(\tau_* - \tau_i) T_*$	2.0	1.7	1.4	1.1	0.84	0.85	1.1	1.5	1.8	2.1
$(\tau_* - \tau_f) T_*$	0.00	0.05	0.11	0.19	0.24	0.24	0.20	0.11	0.04	0.00
$\frac{\mathcal{P}_\perp(\tau_f) - \mathcal{P}_\parallel(\tau_f)}{\mathcal{E}(\tau_f)}$	0.06	-0.03	-0.22	-0.56	-1.1	1.6	0.91	0.47	0.24	0.13

$\tau_*$  = hydro onset time

$T_*$  = initial hydro temperature

$c \rightarrow \infty$ :  $\tau_* \rightarrow \tau_f$ ,  $T_* \rightarrow \infty$ , “instantaneous” relaxation to local equilibrium

$c \rightarrow 0$ :  $\tau_* \sim \Delta/\sqrt{c}$ ,  $\tau_* T_* \sim O(1)$

$\Lambda\tau_* \geq 0.9$  always  $\Rightarrow$  limit of validity of hydro controlled by relaxation of non-hydro modes, not by growth of higher-order viscous terms

# Open questions

- Sensitivity to choice of boundary time dependence?
  - wider range of amplitudes
  - periodic forcing
- Precise connection between entropy & apparent horizon area?
  - ambiguities in definition of non-equilibrium entropy
  - foliation dependence of apparent horizon area
- Feasibility of evolving anisotropic & inhomogeneous geometries?
  - non-boost invariant colliding shocks
  - finite expanding fluids
  - turbulent driven systems
- Relevance for heavy ion collisions?