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A flavor superconductor from string theory

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## Motivation

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Holographic Superconductor from charged scalar in Einstein-Maxwell gravity

Gubser; Hartnoll, Herzog, Horowitz

cf. Chris Herzog's talk

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p-wave superconductor (current dual to gauge field condensing)

Gubser, Pufu

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( $AdS_4$  examples)

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1) for which the field theory is explicitly known ?

2) within ten-dimensional type IIB supergravity?



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A holographic superconductor with field theory in 3+1 dimensions for which

1. the dual field theory is explicitly known
2. there is a qualitative ten-dimensional string theory picture of condensation

Ammon, J.E., Kaminski, Kerner 0810.2316, 0903.1864

p-wave superconductor

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This is achieved in the context of

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$$\phi \rightarrow e^{i\Lambda} \phi e^{-i\Lambda}$$

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Brane probes added on gravity side  $\Rightarrow$  fundamental d.o.f. in the dual field theory

Additional hyperplanes within  $AdS_5 \times S^5$  or deformed version thereof

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$\Rightarrow$  Rich phase structure

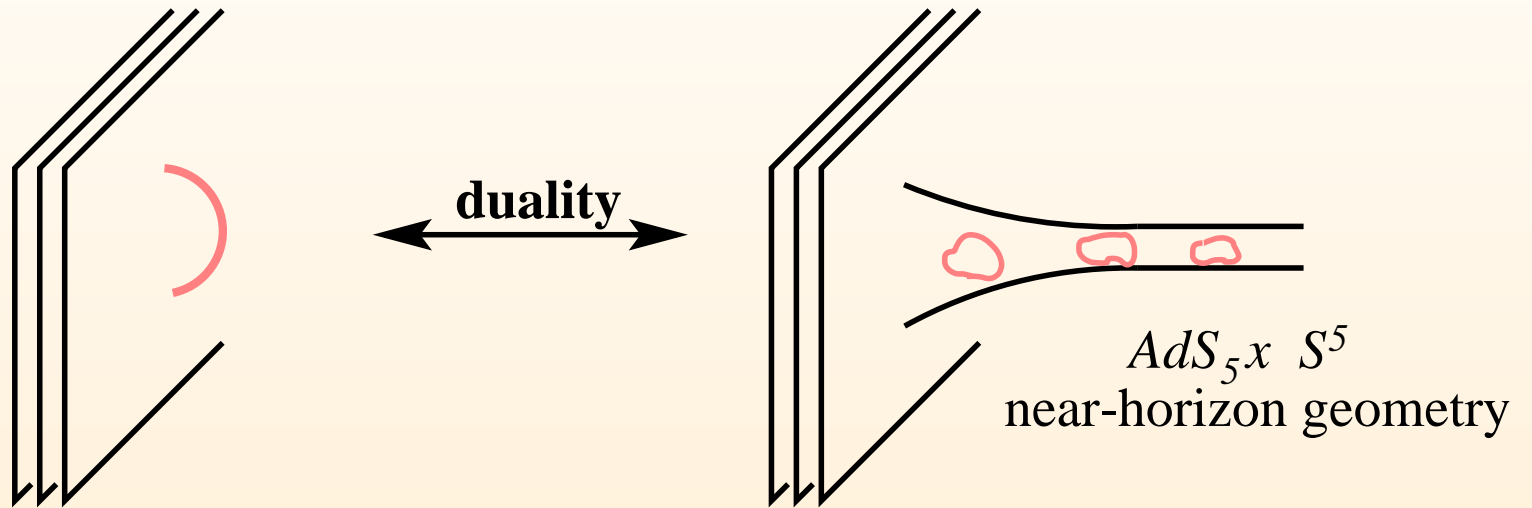
# Outline

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1. Adding Flavor to Gauge/Gravity Duality
2. Holographic Quarks at finite Temperature and Density
3. Superconductivity

# String theory origin of AdS/CFT correspondence

D3 branes in 10d



↓ Low-energy limit

$\mathcal{N} = 4$  SUSY  $SU(N)$  gauge theory in four dimensions  
( $N \rightarrow \infty$ )

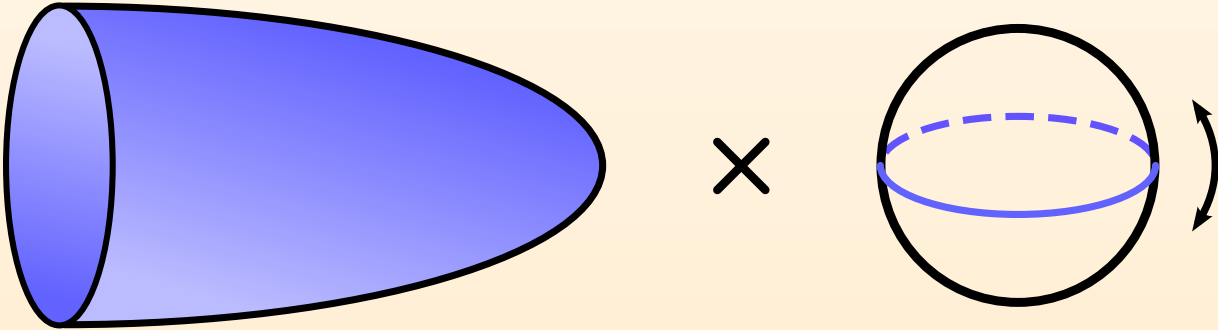
Supergravity on  $AdS_5 \times S^5$

# Quarks (fundamental fields) within the AdS/CFT correspondence

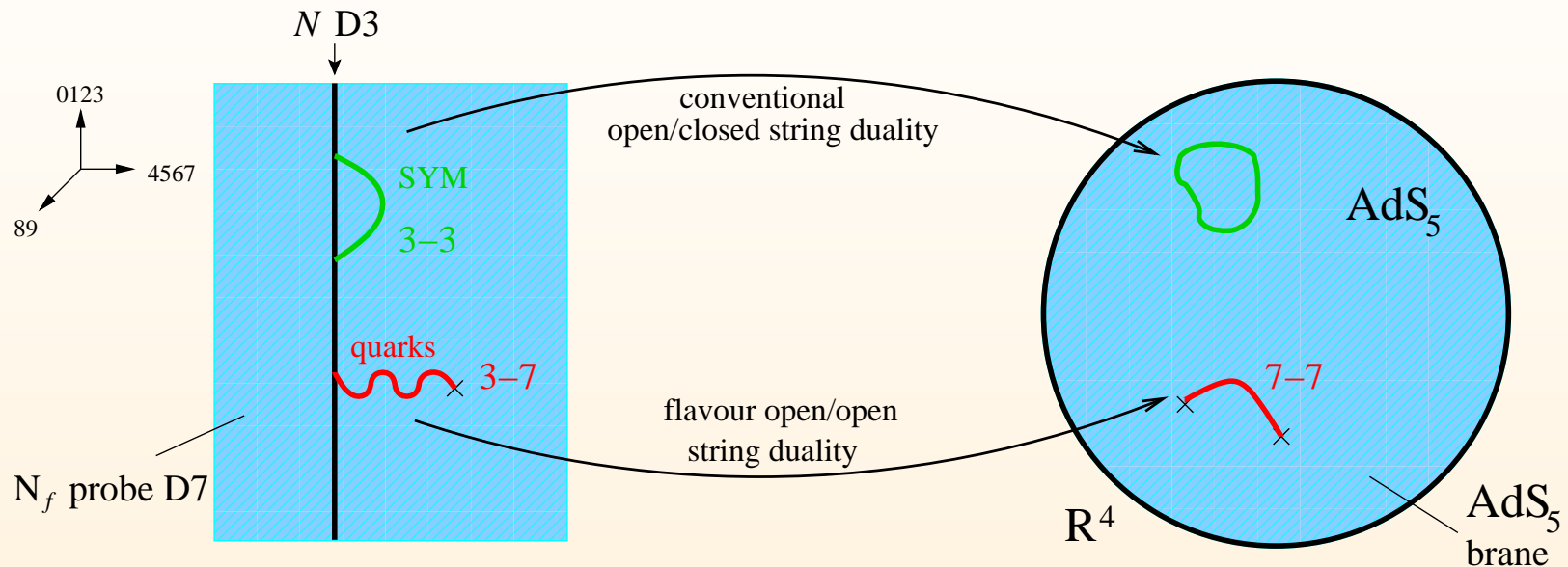
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Adding D7 brane probe:

|    | 0 | 1 | 2 | 3 | 4 | 5 | 6 | 7 | 8 | 9 |
|----|---|---|---|---|---|---|---|---|---|---|
| D3 | X | X | X | X |   |   |   |   |   |   |
| D7 | X | X | X | X | X | X | X | X |   |   |



# Quarks (fundamental fields) from brane probes



$N \rightarrow \infty$  (standard Maldacena limit),  $N_f$  small (probe approximation)

duality acts twice:

$\mathcal{N} = 4$  SU(N) Super Yang-Mills theory

coupled to

$\mathcal{N} = 2$  fundamental hypermultiplet

$\longleftrightarrow$

IIB supergravity on  $AdS_5 \times S^5$

+

Probe brane action on  $AdS_5 \times S^3$

Karch, Katz 2002

Dirac-Born-Infeld action



## Gauge/Gravity Duality with Flavor

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DBI (Dirac-Born-Infeld) action:

$$S_{DBI} = -T_7 \int d^8\xi \operatorname{tr} \sqrt{\det(-P[G] + 2\pi\alpha'F)}$$

Contributions of order  $N_f/N_c$

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Field theory involves fundamental fermions and scalars

# Gauge/Gravity Duality at Finite Temperature

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$\mathcal{N} = 4$  Super Yang-Mills theory at finite temperature is dual to **AdS black hole**

Witten 1998

$$ds^2 = \frac{1}{2} \left( \frac{\varrho}{R} \right)^2 \left( -\frac{f^2}{\tilde{f}} dt^2 + \tilde{f} d\vec{x}^2 \right) + \left( \frac{R}{\varrho} \right)^2 (d\varrho^2 + \varrho^2 d\Omega_5^2)$$

$$f(\varrho) = 1 - \frac{\varrho_H^4}{\varrho^4}, \quad \tilde{f}(\varrho) = 1 + \frac{\varrho_H^4}{\varrho^4}$$

Temperature and horizon related by

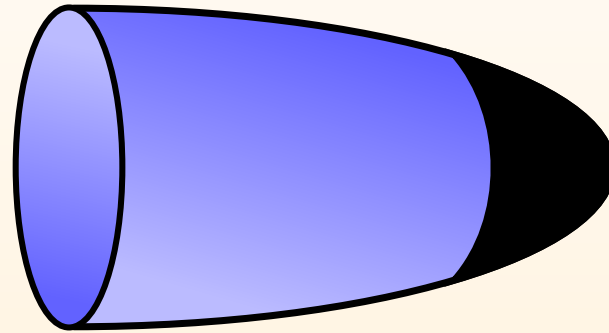
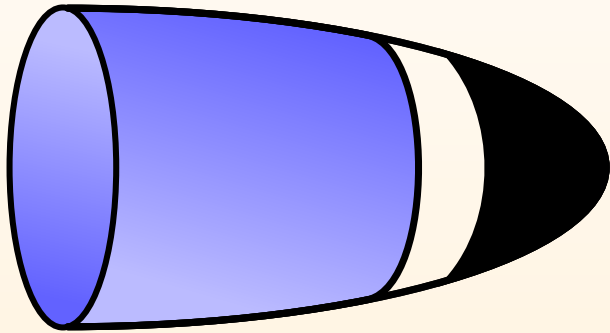
$$T = \frac{\varrho_H}{\pi R^2}$$

$R$ : AdS radius

For  $\varrho_H \rightarrow 0$ , metric of  $AdS_5 \times S^5$  is recovered.

## D7 brane embedding in black hole background

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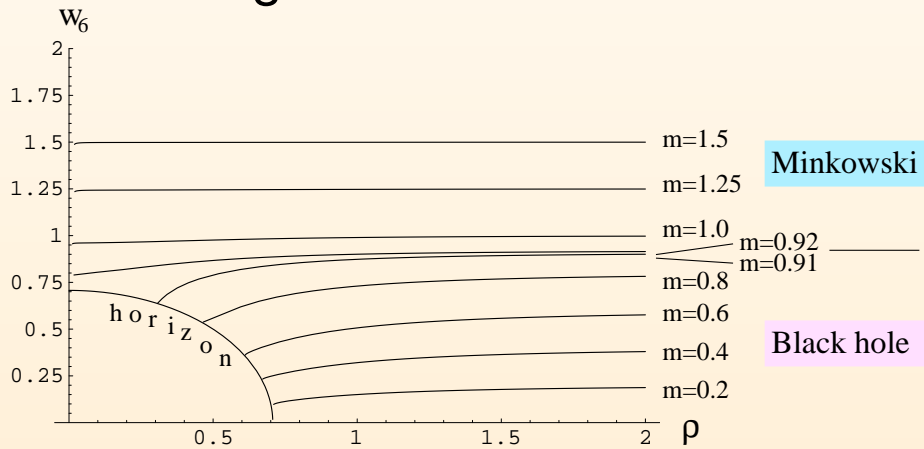
First order phase transition

Babington, J.E., Evans, Guralnik, Kirsch  
Mateos, Myers, Thomson

# D7 brane embedding in black hole background

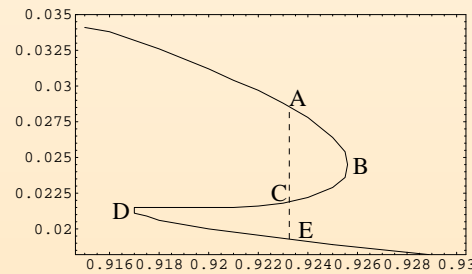
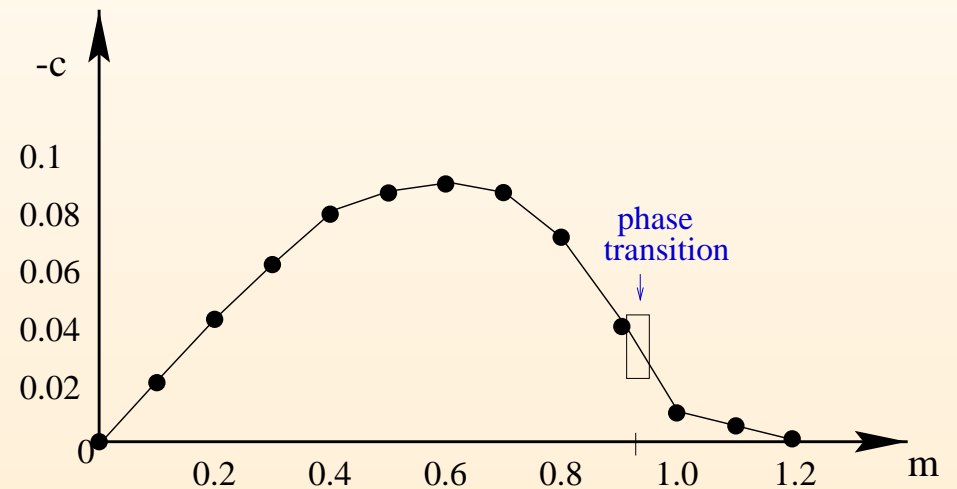
Babington, J.E., Evans, Guralnik, Kirsch 0306018

## Embeddings



Phase transition at  $m_c \approx 0.92$   
(1st order)

Condensate  $c \equiv \langle \bar{\psi}\psi \rangle$  vs. quark mass  $m$   
 $m$  in units of  $T$



Kirsch 2004

## Masses and decay widths of mesons - Spectral functions

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Standard procedure in D3/D7:

Meson masses calculated from linearized fluctuations of D7 embedding

Fluctuations:  $\delta w(x, \rho) = f(\rho) e^{i(\vec{k} \cdot \vec{x} - \omega t)}$ ,  $M^2 = -k^2$

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Identify mesons with resonances in spectral function

Spectral function determined by poles of retarded Green function

Quasinormal modes



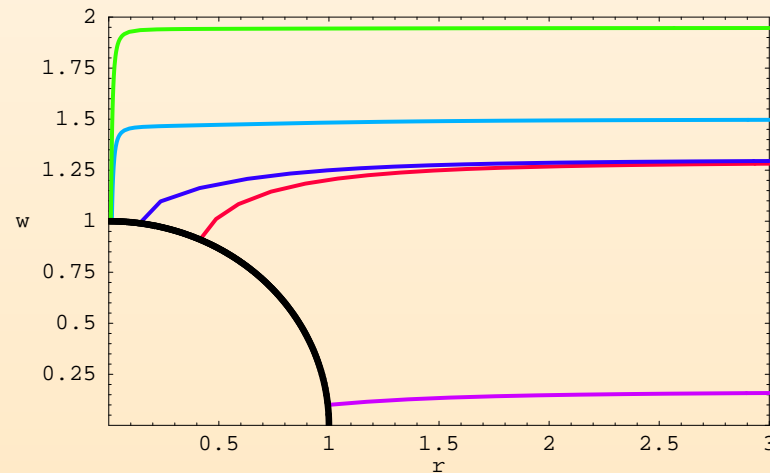
# Finite $U(1)$ baryon density

Mateos, Myers, Matsuura et al

Baryon density  $n_B$  and  $U(1)$  chemical potential  $\mu$   
from VEV for gauge field time component:

$$\bar{A}_0(\rho) \sim \mu + \frac{\tilde{d}}{\rho^2}, \quad \tilde{d} = \frac{2^{5/2}}{N_f \sqrt{\lambda} T^3} n_B$$

At finite baryon density, all embeddings are black hole embeddings

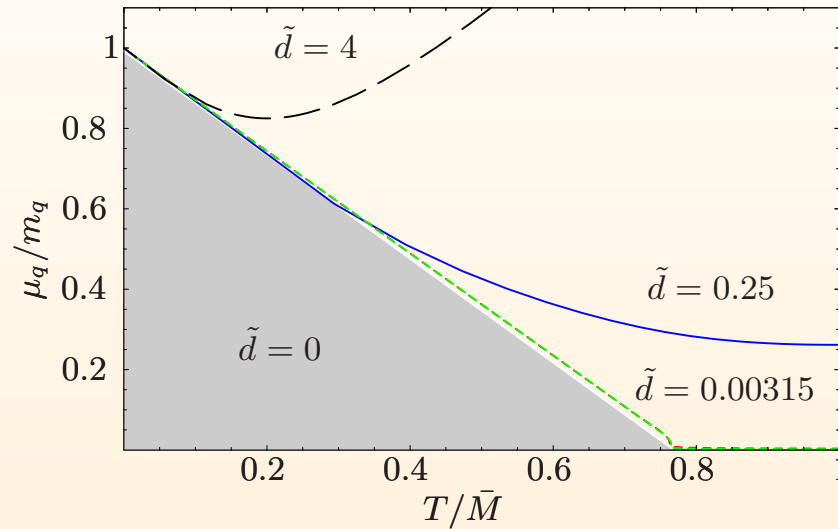


# Phase diagram with finite $U(1)$ baryon density

Phase diagram:

grey region:  $n_B = 0$

white region:  $n_B \neq 0$



Sin, Yogendran et al; Mateos, Myers et al; Karch, O'Bannon; ...

## Isospin chemical potential and density

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- Embed two coincident D7-branes into AdS-Schwarzschild  
gauge fields  $A_\mu = A_\mu^a \sigma^a \in u(2) = u(1)_B \oplus su(2)_I$
- Dynamics of Flavour degrees is described by non-abelian DBI action
- Finite isospin density:  $A_0^3 \neq 0 \Rightarrow$  Explicit breaking to  $u(1)_3$

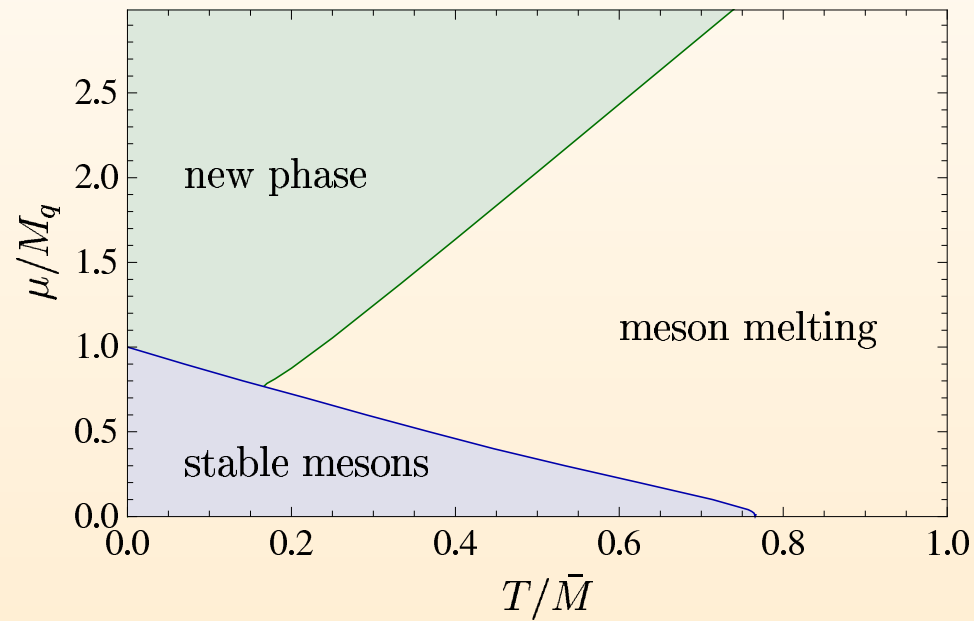
Field theory described:

$\mathcal{N} = 4$  Super Yang-Mills plus two flavors of fundamental matter  
at finite temperature and finite isospin density

# $\rho$ meson condensation

J.E., Kaminski, Kerner, Rust 0807.2663

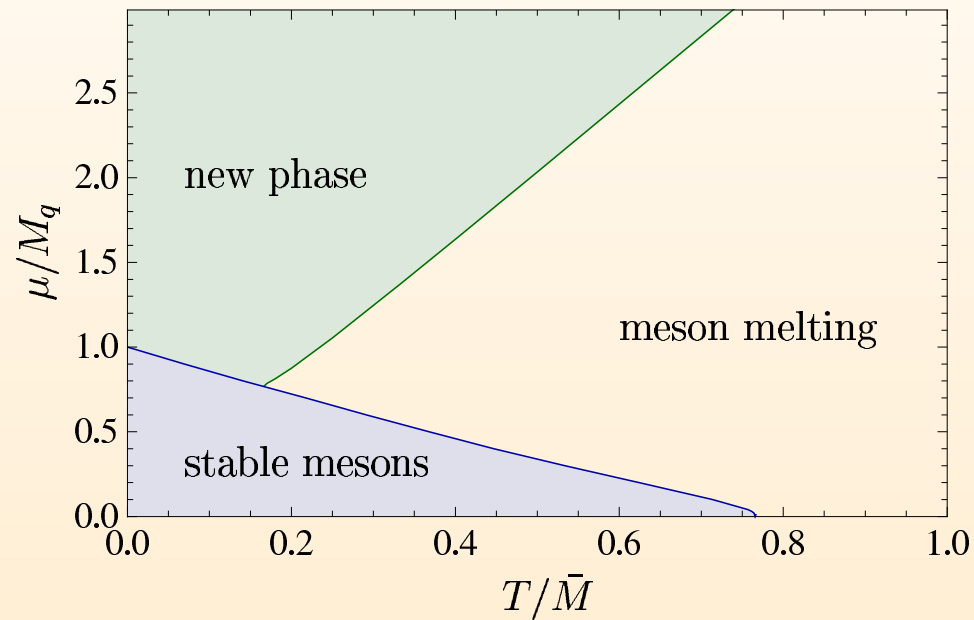
Above a critical isospin density, a new phase forms



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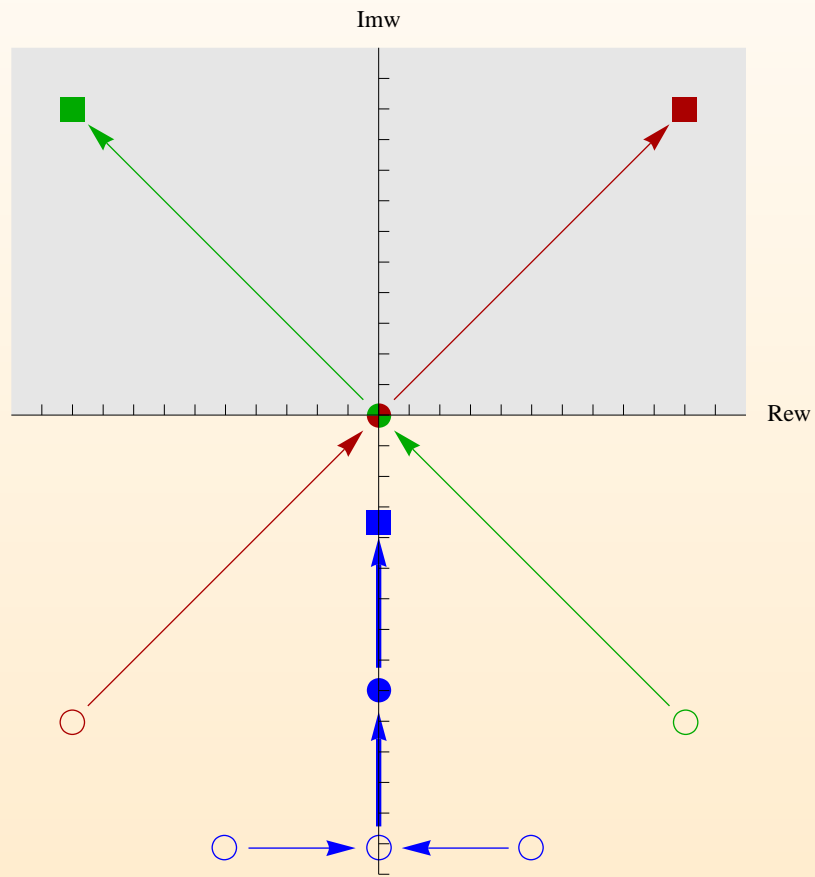
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New phase is unstable

# Quasinormal modes

Instability:



## A new ground state forms

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There is a new solution to the equations of motion

with non-zero vev for  $A_3^1 \sigma^1$  in addition to the non-zero  $A_0^3 \sigma^3$

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$$A_0^3 = \mu - \frac{\tilde{d}_0^3}{2\pi\alpha'} \frac{\rho_H}{\rho^2} + \dots, \quad A_3^1 = -\frac{\tilde{d}_1^3}{2\pi\alpha'} \frac{\rho_H}{\rho^2} + \dots$$



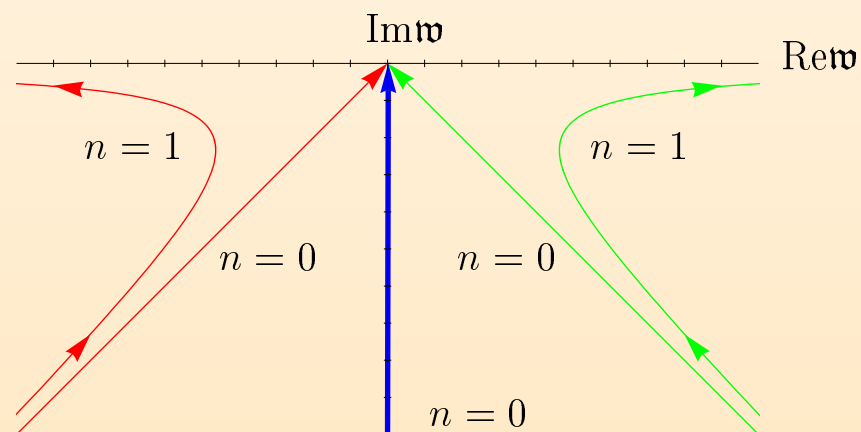
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Pole structure:



# Superconductivity

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Ammon, J.E., Kaminski, Kerner 0810.2316, 0903.1864

The new ground state has properties known from superconductors:

- infinite DC conductivity, gap in the AC conductivity
- second order phase transition, critical exponent of  $1/2$  (mean field)
- a remnant of the Meissner–Ochsenfeld effect

# Superconductivity

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Order parameter  $\tilde{d}_3^1 \propto \langle \bar{\psi}_u \gamma_3 \psi_d + \bar{\psi}_d \gamma_3 \psi_u + \text{bosons} \rangle \neq 0$

Dual to  $A_3^1 \sigma^1$  in gravity theory

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Spontaneous breaking of (global)  $U(1)_3$

Flavor superfluid

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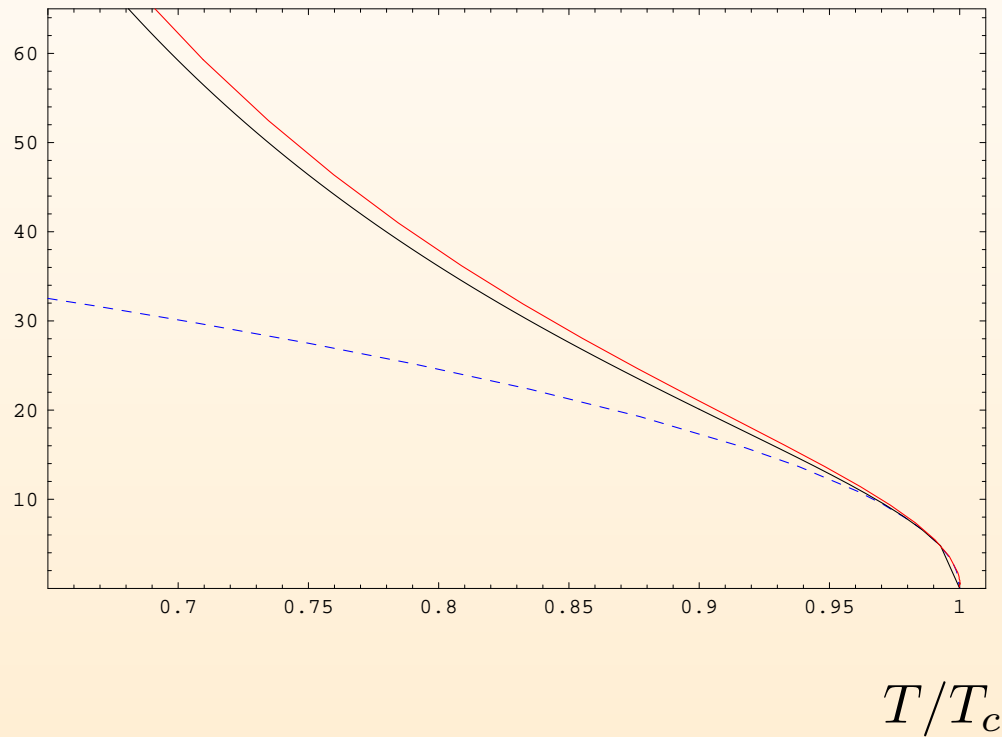
Spontaneous breaking of (global)  $U(1)_3$

Flavor superfluid

$U(1)_3$  may be weakly gauged

# Order parameter: p wave condensate

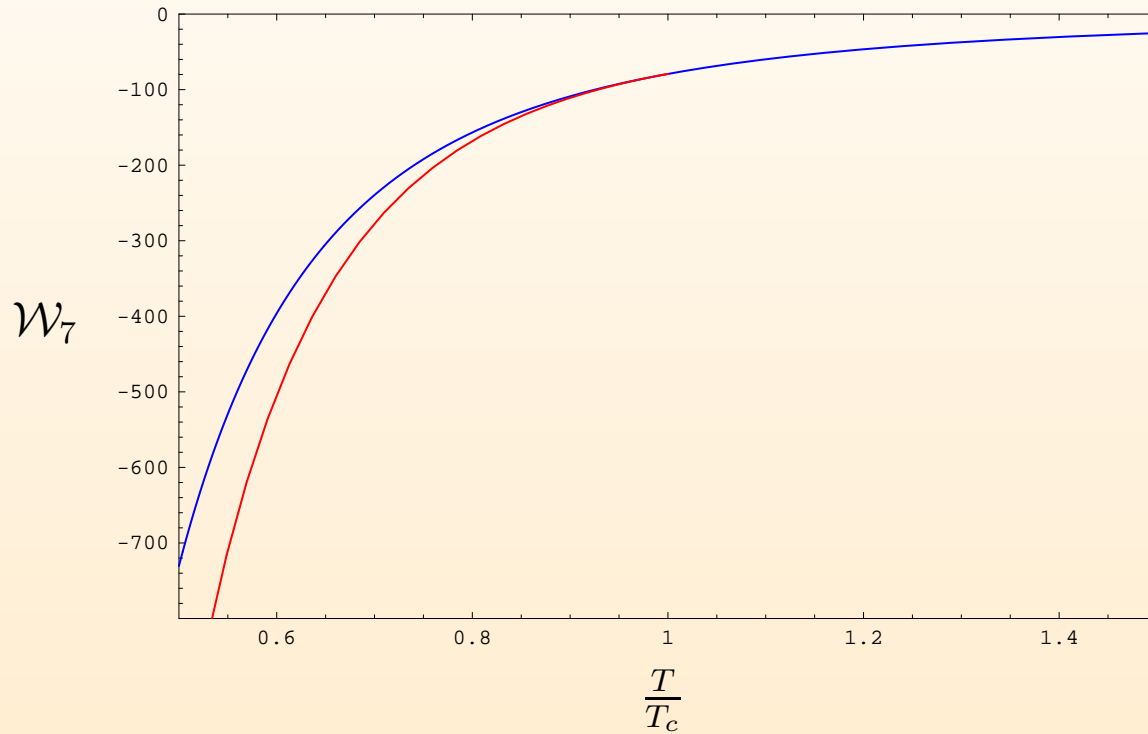
$-\tilde{d}_3^1$



Red: Vanishing quark mass;      Black: Finite quark mass,  $\mu/M_q = 3$

Blue: Fit displaying critical exponent 1/2

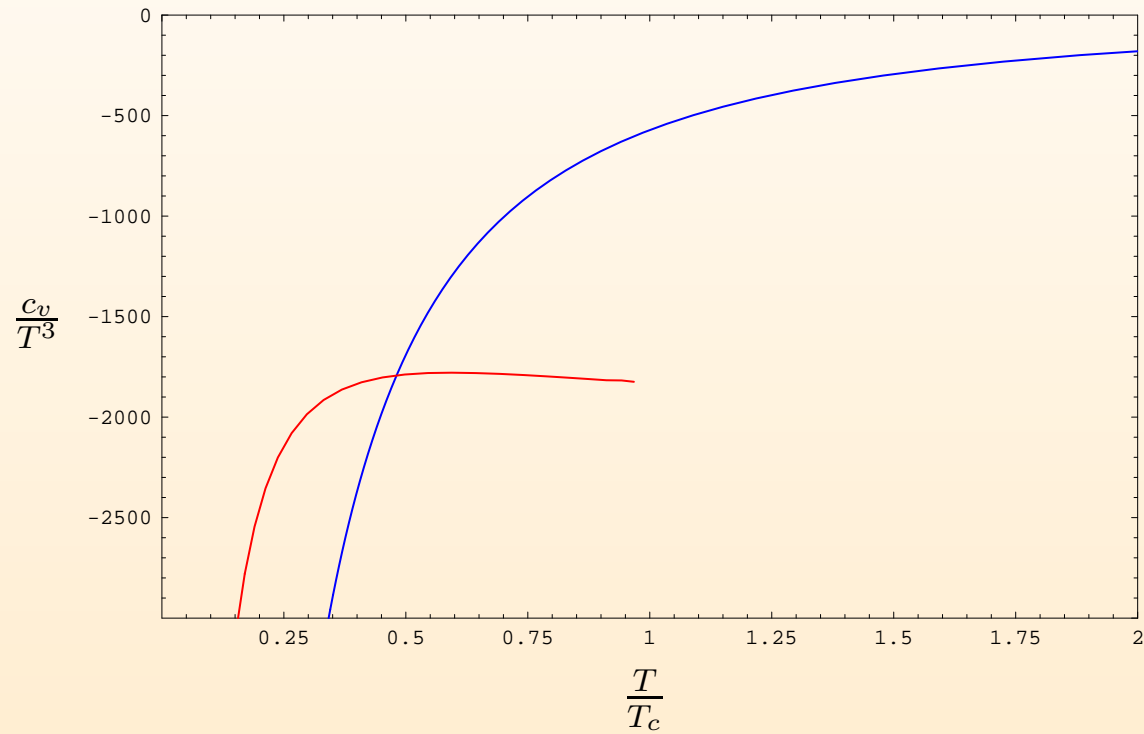
## Flavor contribution to Grand potential vs. temperature



# Heat Capacity

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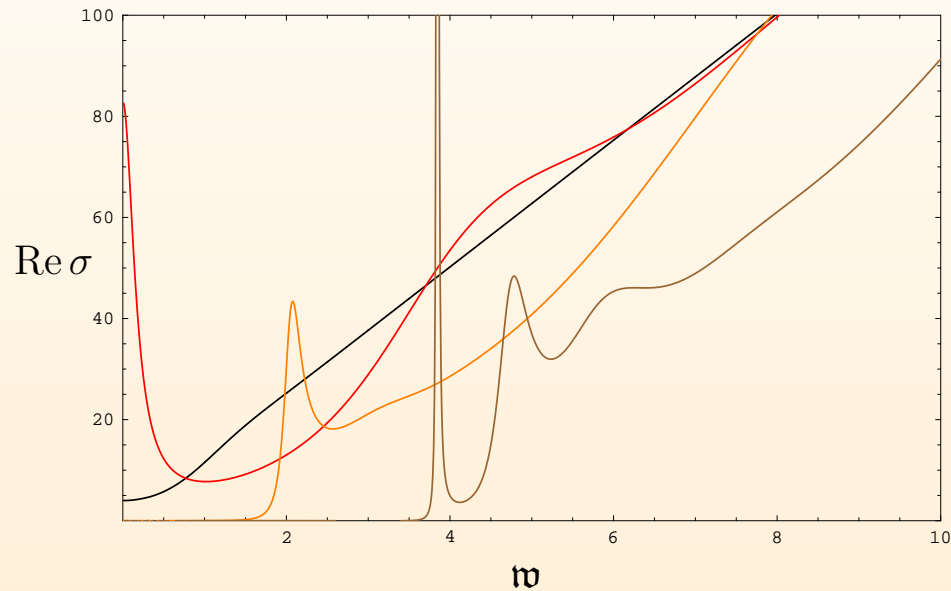
## Flavor contribution to heat capacity





# Conductivity

## Frequency-dependent conductivity

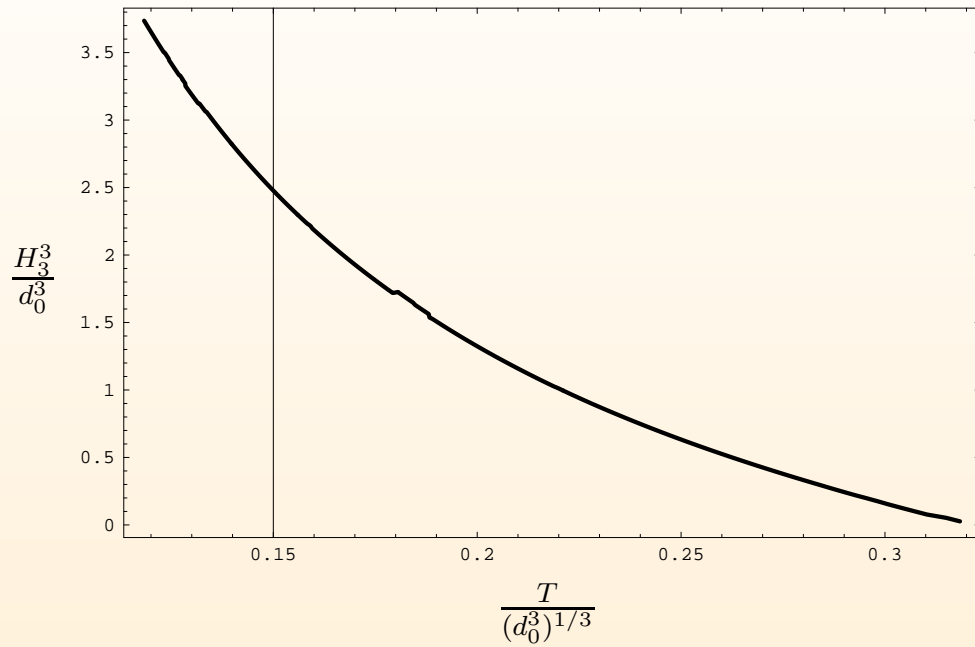


$$\nu = \omega / (2\pi T)$$

$T/T_c$ : Black:  $\infty$ , Red: 1, Orange:  $T = 0.5$ , Brown:  $T = 0.28$ .

(Vanishing quark mass)

# Meissner effect



Lower phase: magnetic field and condensate coexist

Upper phase: condensate vanishes

## Non-abelian DBI action

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Evaluation non-trivial in presence of both  $\sigma^0, \sigma^1$

Two evaluation methods:

1) Expansion to fourth order

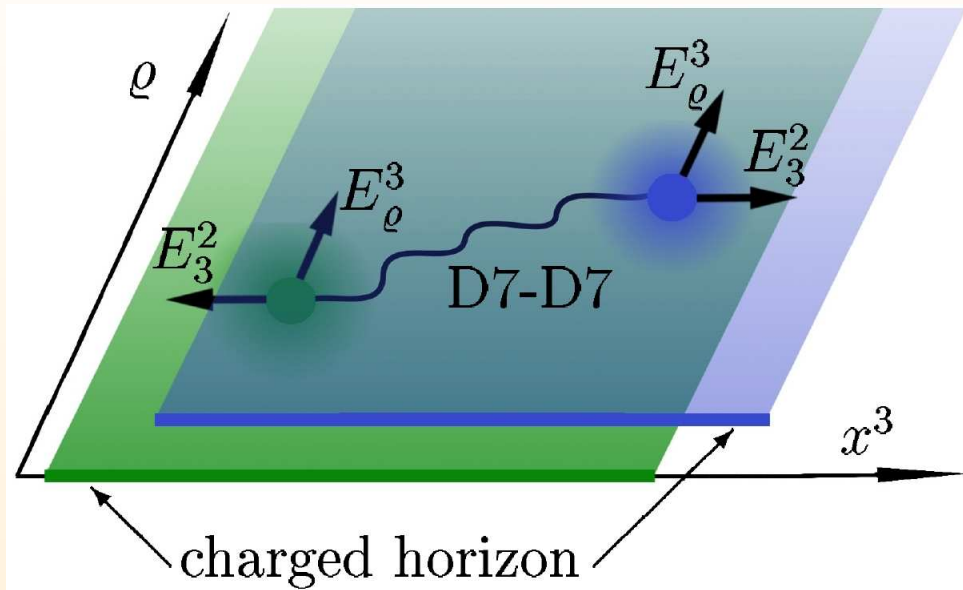
2) Simplification: Omitting commutators of Pauli matrices  
Modified prescription for symmetrized trace

Allows for all-order calculation of the non-abelian DBI

Error of order  $1/N_f$

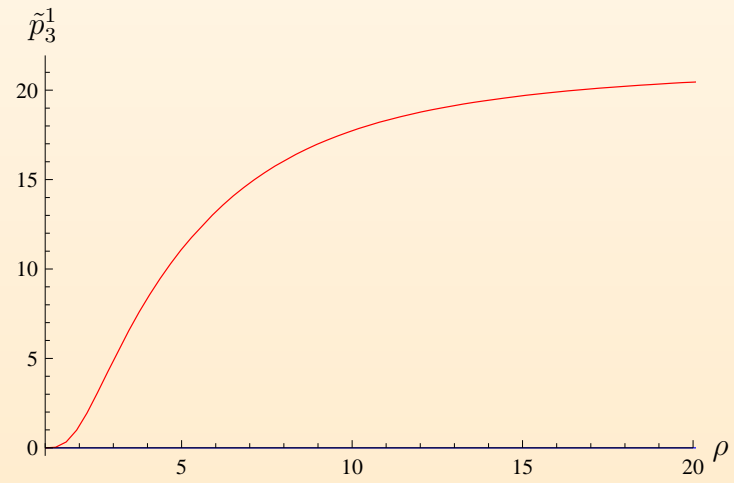
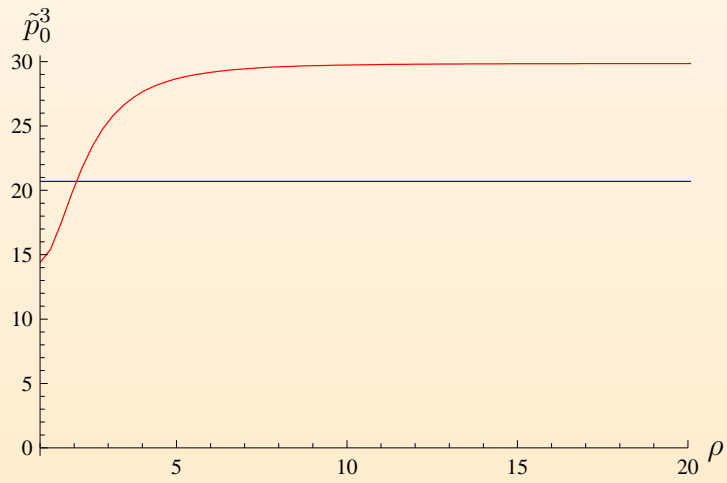
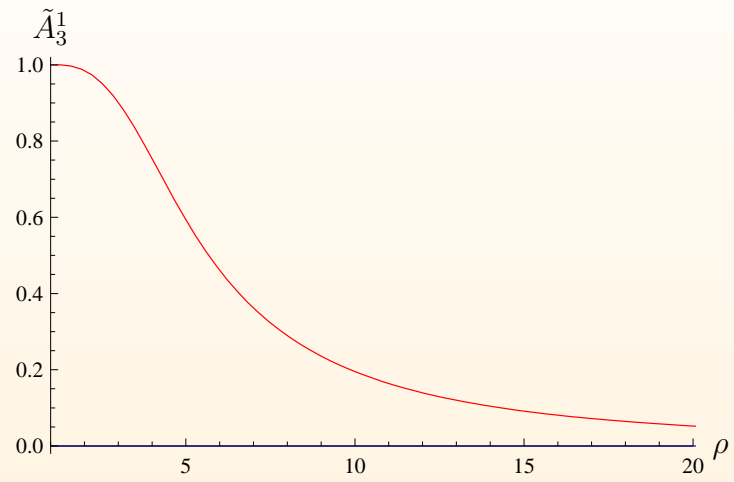
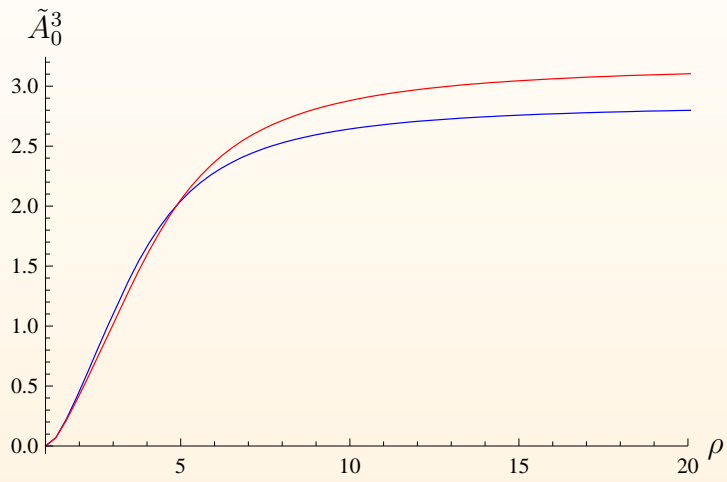
cf. Myers, Constable, Tafjord 1999

## String picture



- Strings stretched between D7 branes and horizon induce a charge near the horizon
- System unstable above a critical charge density
- Horizon strings recombine to  $D7 - D7$  strings
- $D7 - D7$  strings propagate into the bulk, balancing flavorelectric and gravitational forces
- $D7 - D7$  strings distribute isospin charge into the bulk  $\rightarrow$  superconducting condensate

# Charge distributions



## Conclusion and Outlook

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- A holographic superconductor for which the field theory is explicitly known
- First explicit example of superconductivity (superfluidity) from 10d action
- Embedding of two coincident D7 branes  $\Rightarrow$  Finite isospin density

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- Outlook: Fermions
  
- Outlook: Space-time dependent solutions  
Spin density waves (w. E. Caceres)